

# How can birds and bacteria move together without a leader? An introduction to collective motion in biology

Fernando Peruani

*In collaboration with:*

***M. Bär, H. Chaté, A. Deutsch, F. Ginelli, V. Jakovlevic, L. Søgaard-Andersen, and J. Starruß***

**Summer Solstice Conference - Nancy – 2010**



**Motivation: Examples of collective motion in biology**

**Collective motion in a simple model**

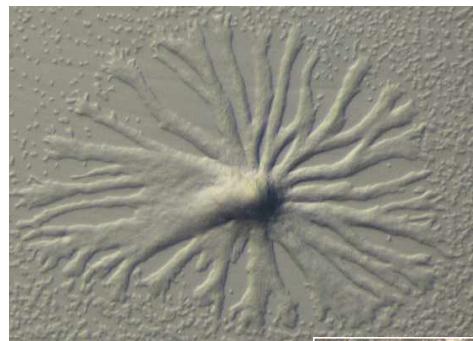
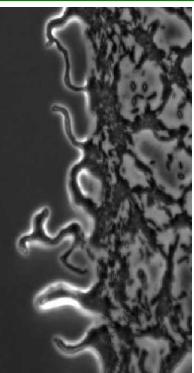
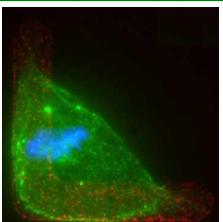
**A specific example: collective motion in myxobacteria**

**Symmetries!**

**Collective motion on the lattice**

**Summary**

# Motivation: Examples of collective motion in biology



Molecular  
motors

bacteria  
 $\sim 10^{-6}$  m

eukaryote cells



social insects

$\sim 10^{-2}$  m

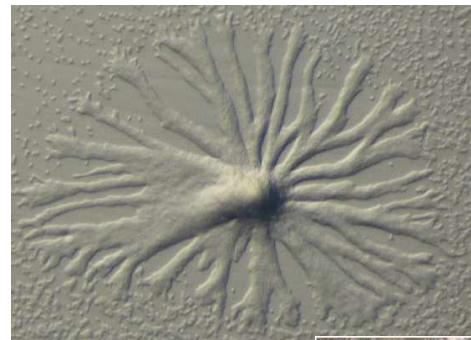
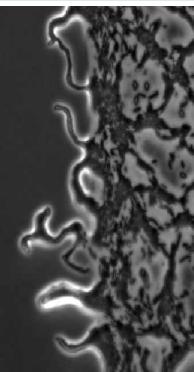
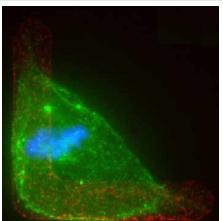


bird and fish  
 $\sim 10^{-1}$  m



Mammals  
 $\sim 10^0$  m

# Motivation: Examples of collective motion in biology



Molecular  
motors

bacteria  
 $\sim 10^{-6}$  m

eukaryote cells

social insects

$\sim 10^2$  m

Large-scale patterns of millions of individuals emerge without a central control system!

Coherent moving structures that emerge from local rules and short range interactions – i.e., without global knowledge of the system!

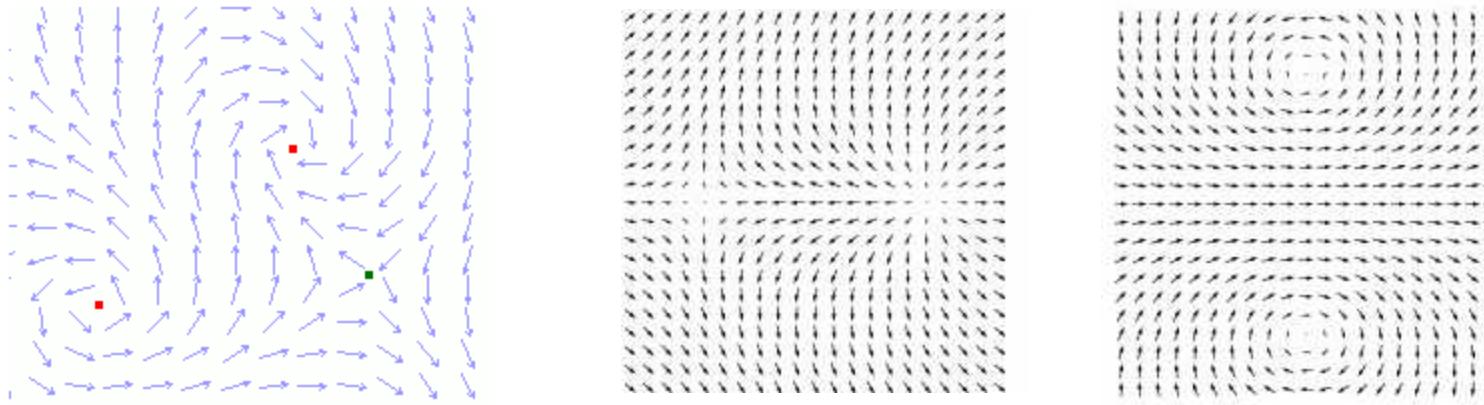


# Motivation: collective motion as a theoretical challenge

Imperfect flow of information leads to defects, and defects set the limit to the size/scale of the patterns we can observe

Some classical results from statistical mechanics:

- The Kosterlitz-Thouless transition:  $E_{ij} = -J(s_i s_j) = -J \cos(\phi_i - \phi_j)$



- The Mermin-Wagner theorem:

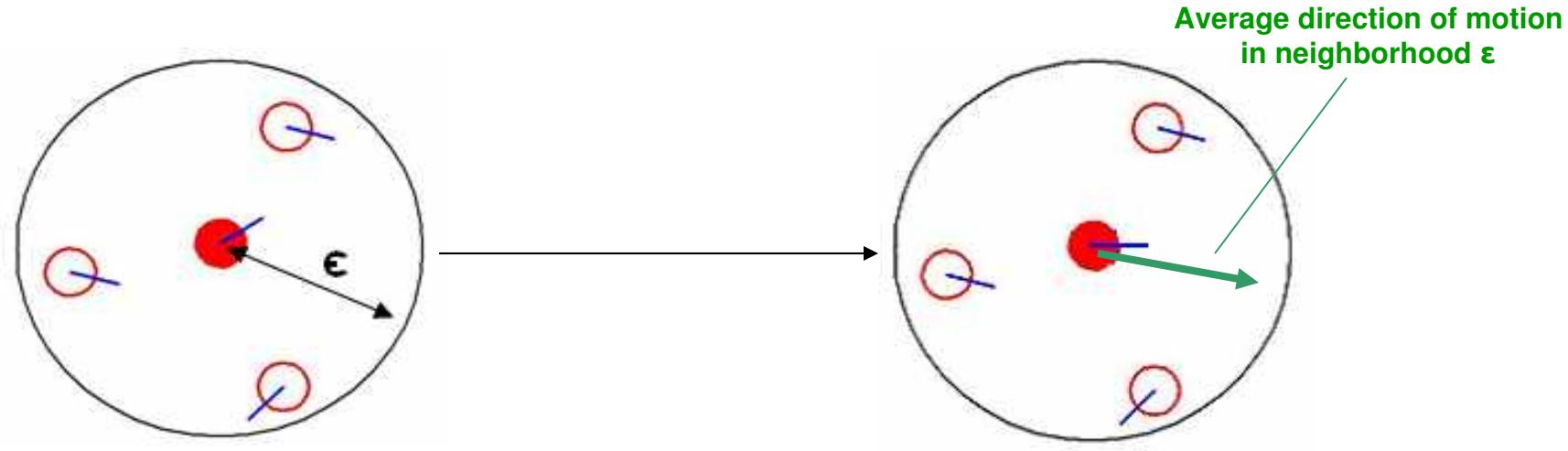
In equilibrium systems with SU2 symmetry, long-range order out of short-range interactions cannot emerge in 1D or 2D!

[Ref.: Peruani, Nicola, Morelli, <http://arxiv.org/abs/1003.4253> (2010)]

# Collective motion in a simple model

## The Vicsek model

[T. Vicsek *et al.*, Phys. Rev. Lett. 75, 1226 (1995)]



Time  $t$

Time  $t+1$

Motion in space

$$\mathbf{x}_i(t + 1) = \mathbf{x}_i(t) + \mathbf{v}_i(t)\Delta t$$

Change of the direction of motion

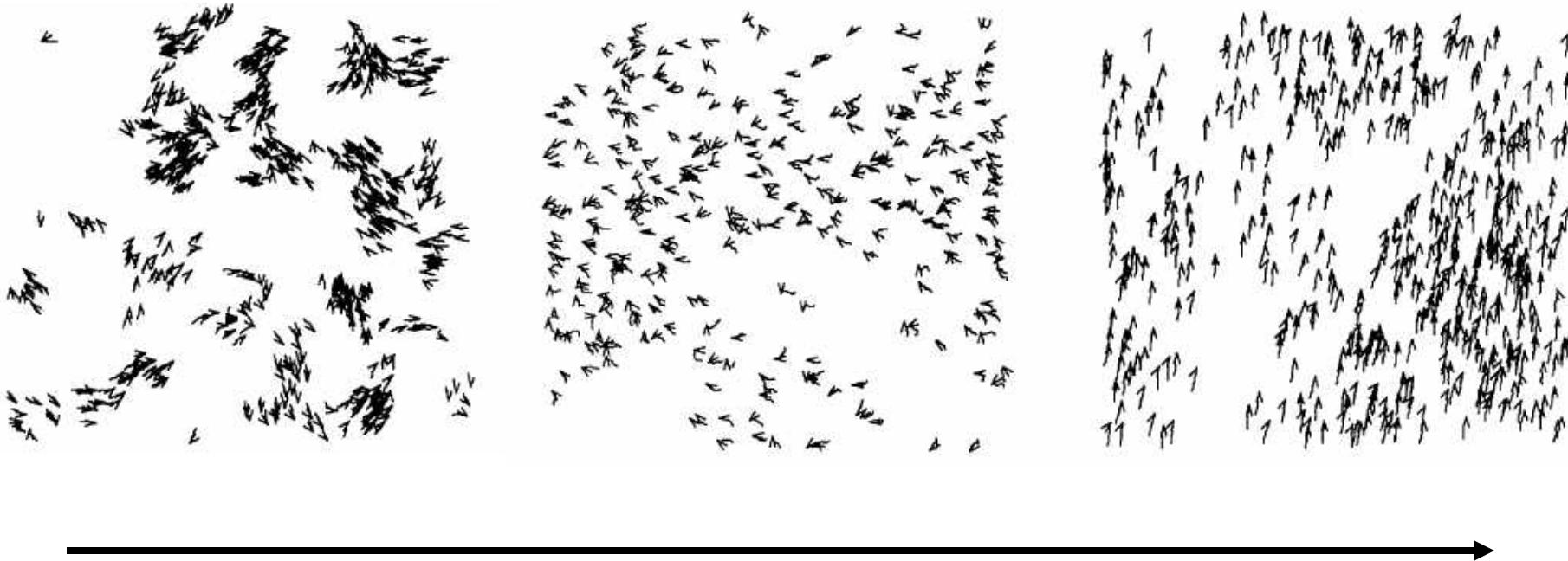
$$\theta_i(t + 1) = \langle\theta_i(t)\rangle_r + \Delta\theta$$

Average direction of motion at time  $t$

Angular noise !!!

# Collective motion in a simple model

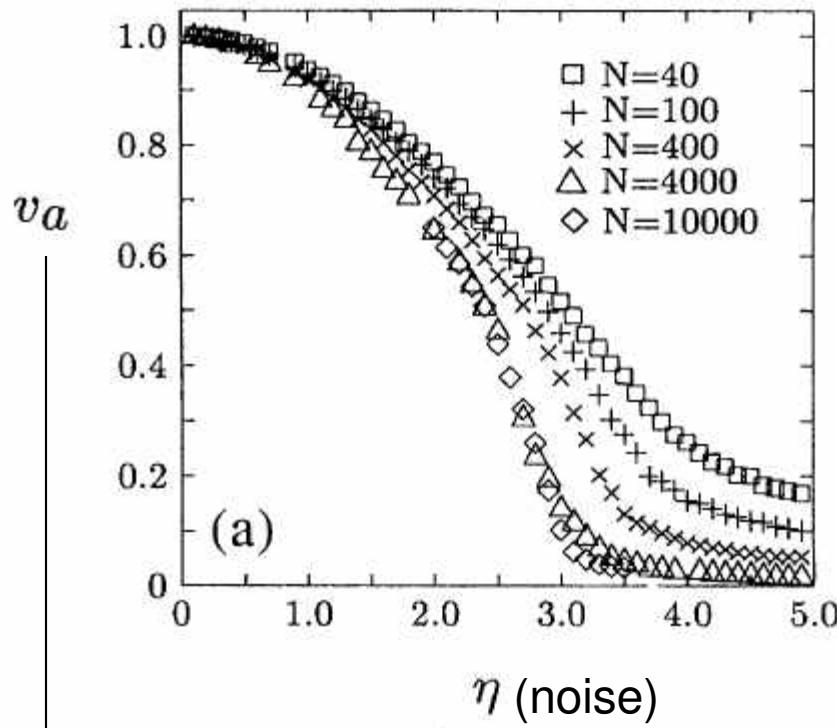
## The Vicsek model



Decreasing noise values

# Collective motion in a simple model

## The Vicsek model

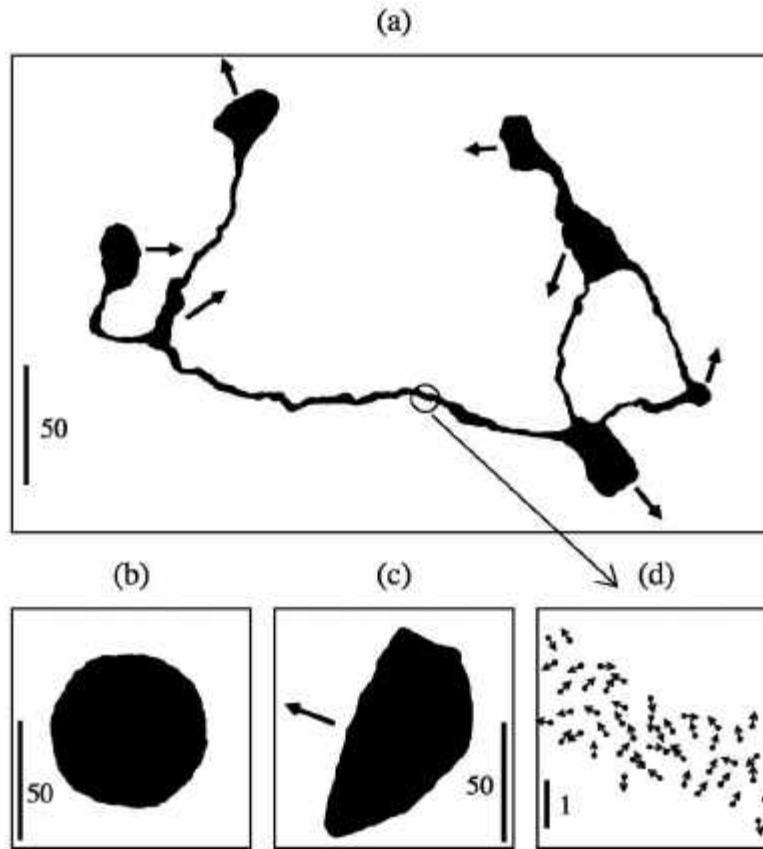


$$v_a = \frac{1}{N} \left| \sum_{i=1}^N \mathbf{v}_i \right|$$

# Collective motion in a simple model

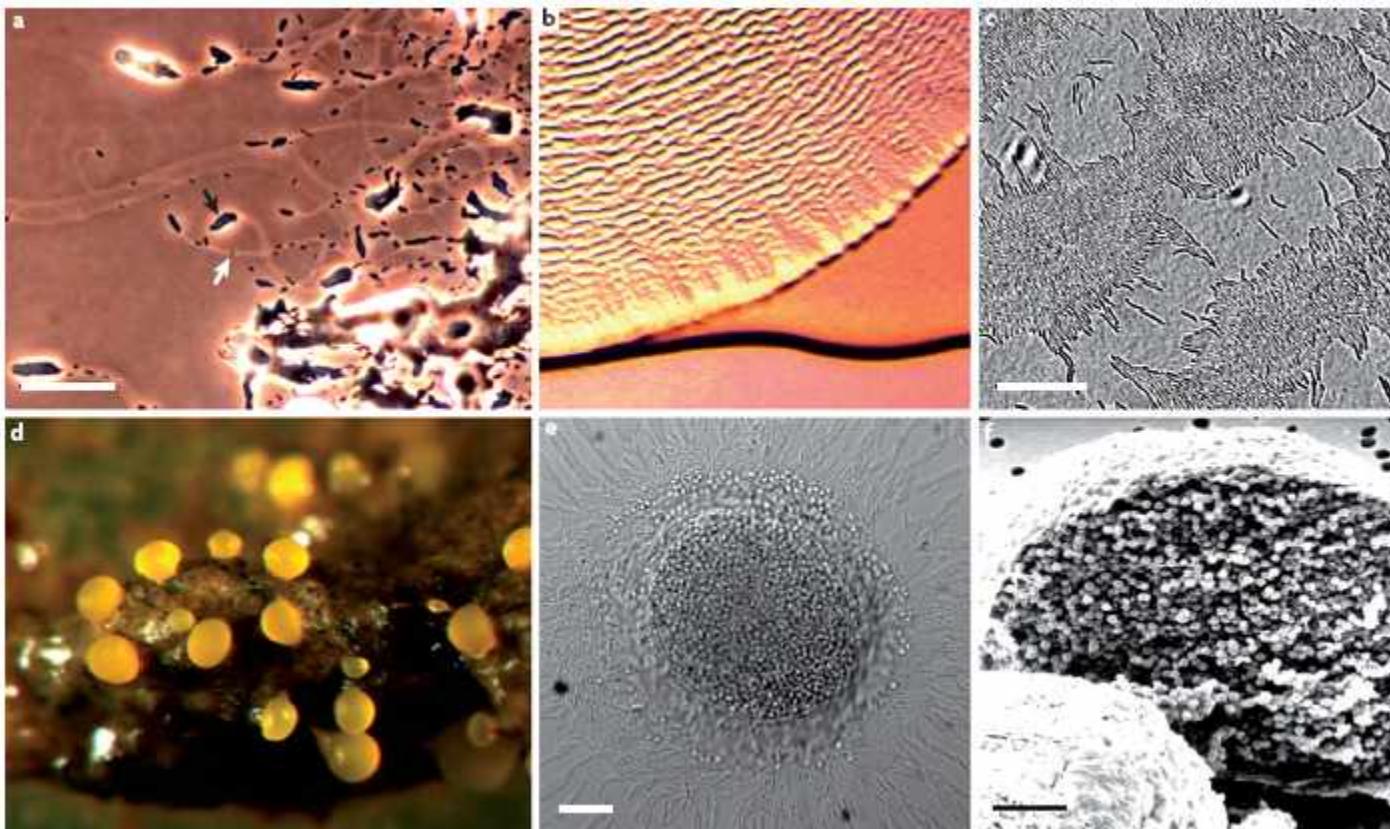
## Need of cohesion to sustain a flock in open space

[G. Grégoire and H. Chaté, Phys. Rev. Lett. 92, 025702 (2004)]



$$\theta_j^{t+1} = \arg \left[ \alpha \sum_{k \sim j} e^{i\theta_k^t} + \beta \sum_{k \sim j} f_{jk}^t e^{i\theta_{jk}^t} + \eta n_j^t e^{i\xi_j^t} \right], \quad \text{where,} \quad f_{jk} = \begin{cases} -\infty & \text{if } r_{jk} < r_c \\ \frac{1}{4} \frac{r_{jk} - r_e}{r_a - r_e} & \text{if } r_c < r_{jk} < r_a \\ 1 & \text{if } r_a < r_{jk} < r_0 \end{cases}$$

# A specific example: collective motion in myxobacteria

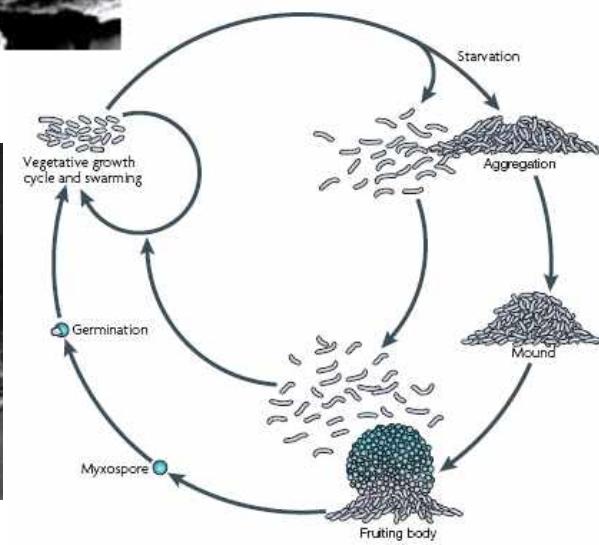
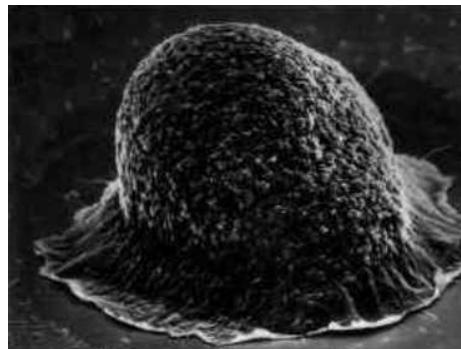
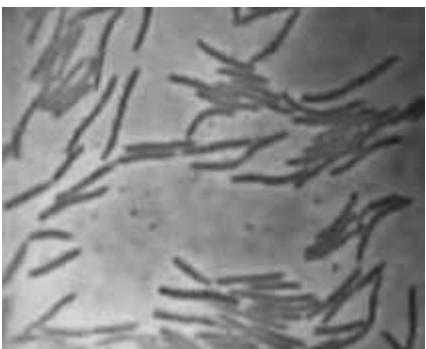


**Myxococcus  
xanthus**

Zusman 2007

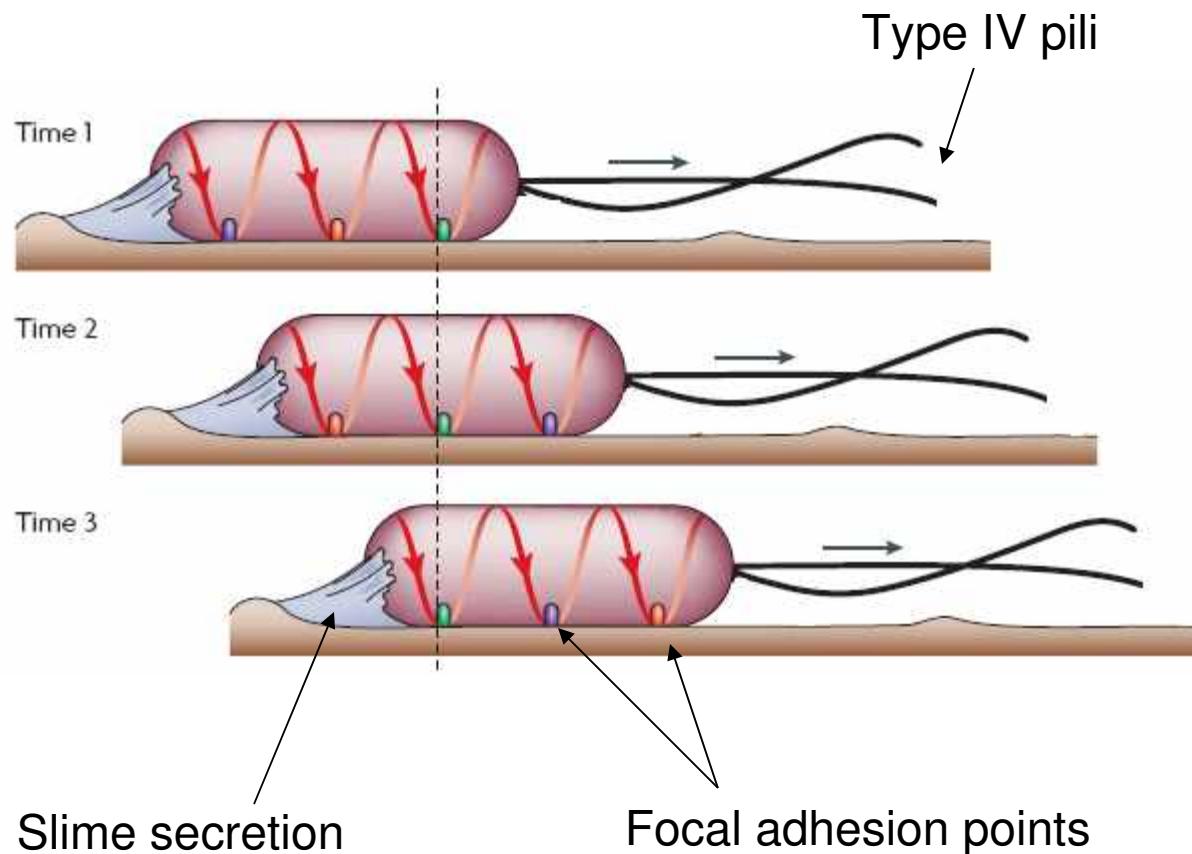
Reichenbach  
1965

F.Peruani



# A specific example: collective motion in myxobacteria

- Motility engines in *M. xanthus*:



Pelling 05

Myxobacteria (speed = 0.025 to 0.1  $\mu\text{m}/\text{s}$ )

Cyanobacteria (speed = 10  $\mu\text{m}/\text{s}$ )

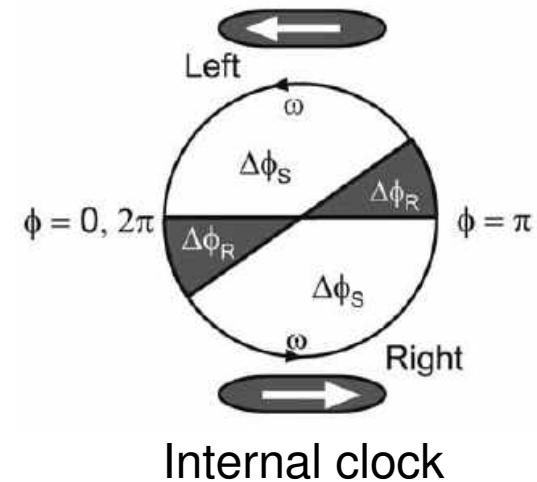
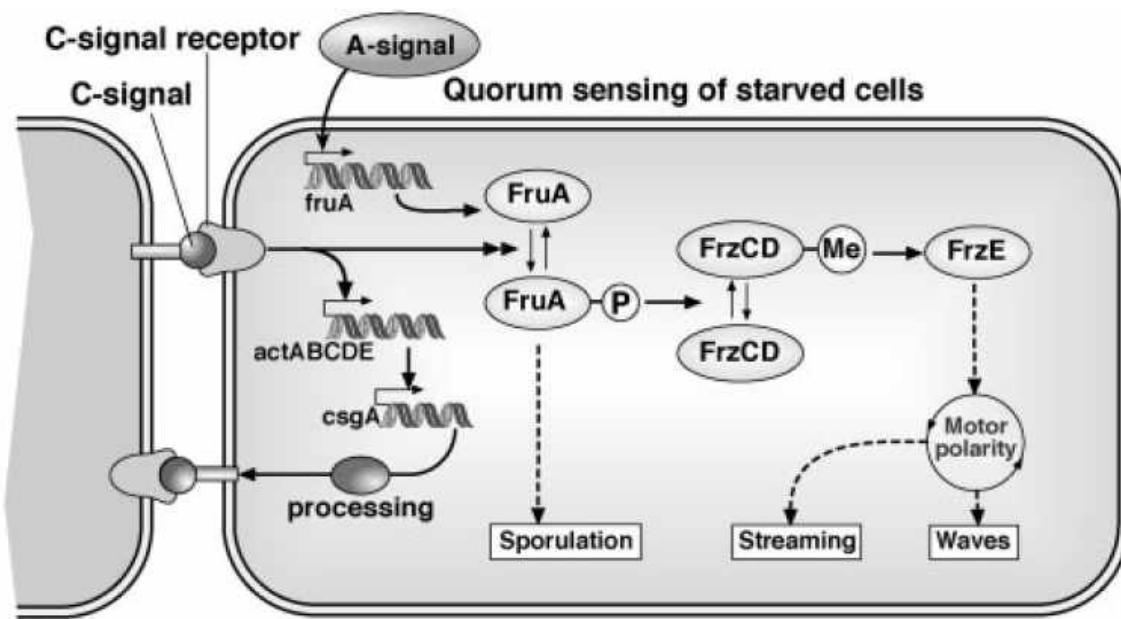
Cytophaga-Flavobacterium (speed = 2 to 4  $\mu\text{m}/\text{s}$ )

# A specific example: collective motion in myxobacteria

- How do *M. xanthus* cells communicate?

- A quorum sensing diffusive mechanism to trigger the life cycle.
- There is no evidence of a guiding chemotactic signals involved in collective motion.
- Cells exchange C-signal which controls cell reversal (it requires cell-cell contact).

- Cell reversal and C-signal:

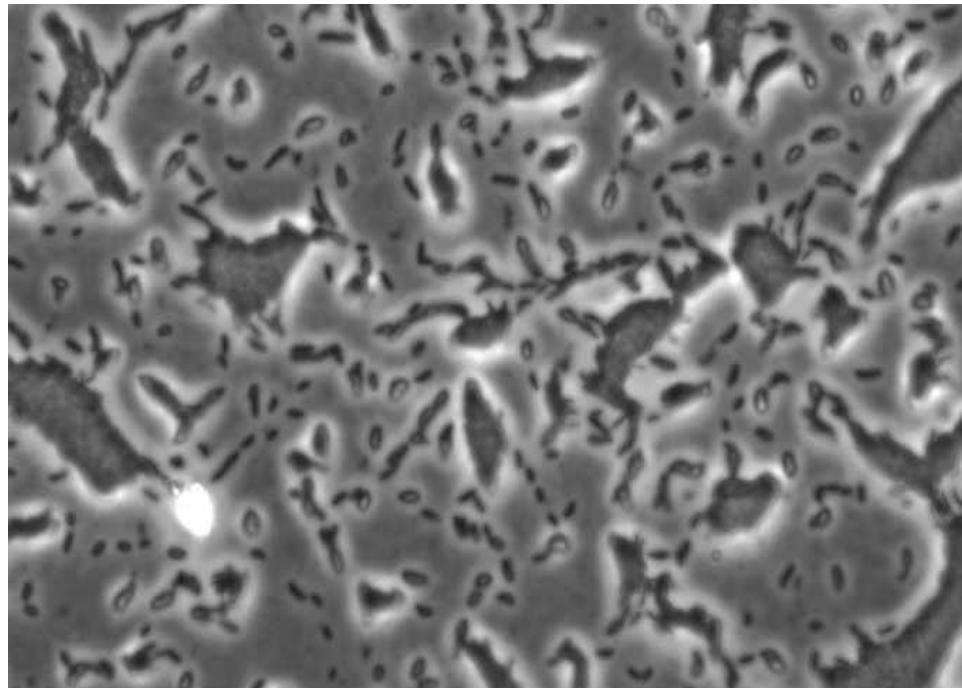
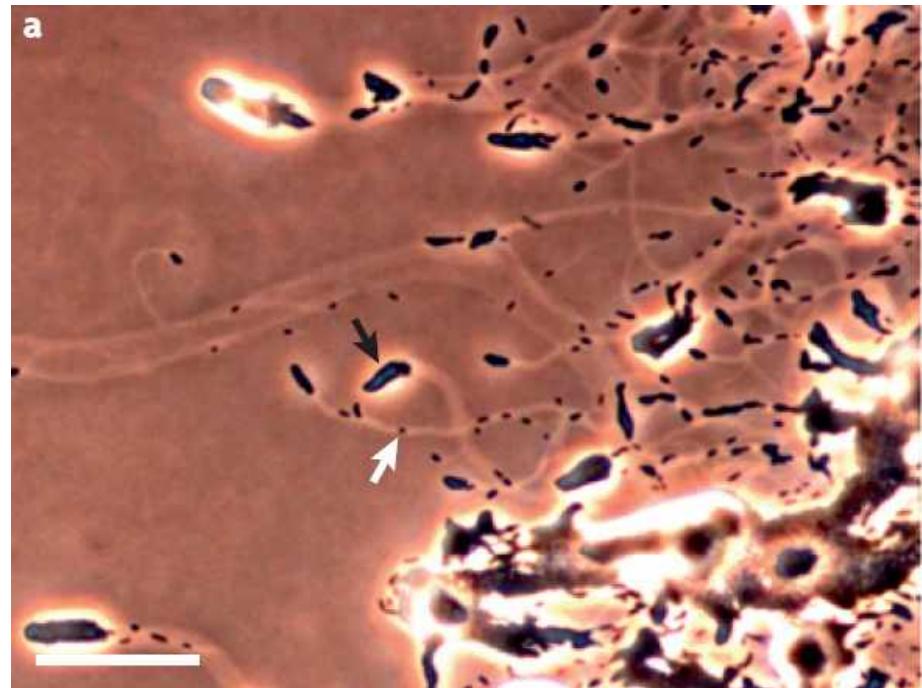


Igoshin & Oster 2003

F.Peruani

# A specific example: collective motion in myxobacteria

Which mechanism is used by the cells to coordinate their motion?



(Collective motion and clustering in the wild type during the vegetative growth)

- Is there a hidden guiding chemotactic signal?
- Can slime trail following cause these effects?
- Is there a cell-density sensing mechanism that controls cell speed causing of these effects?
- What is the minimal mechanism that can produce these effects?

# A specific example: collective motion in myxobacteria

**Self-propulsion of bacteria + elongated shape  
= collective behavior ?**



**What macroscopic effects can we expect  
in a system of self-propelled rods ?**



# A specific example: collective motion in myxobacteria

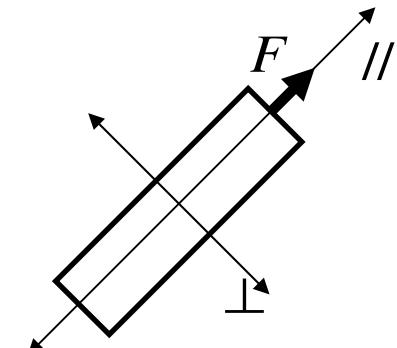
## A simple model for self-propelled rods

[F. Peruani, A. Deutsch, and M. Bär, Phys. Rev. E 74, 030904 (2006)]

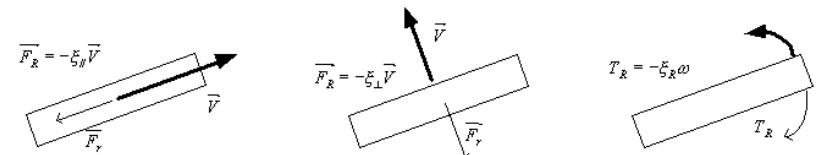
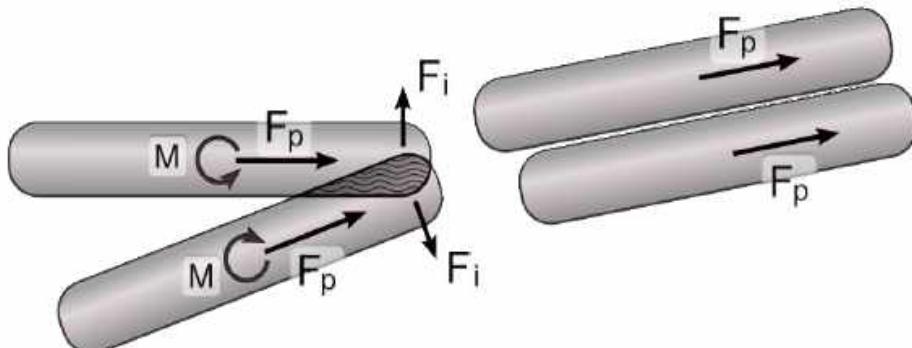
We consider the over-damped situation in which we have:

$$\left\{ \begin{array}{l} (\nu_{\parallel}, \nu_{\perp}) = \left( \frac{1}{\zeta_{\parallel}} (R_{\parallel}(t) + F + F_{Inter_{\parallel}}), \frac{1}{\zeta_{\perp}} (R_{\perp}(t) + F_{Inter_{\perp}}) \right) \\ \frac{d\theta}{dt} = \frac{1}{\zeta_R} (\tilde{R}(t) + M_{Inter}) \end{array} \right.$$

Self-Propelling force  
Interactions



$R_{\parallel}, R_{\perp}, \tilde{R}$  are white noises!

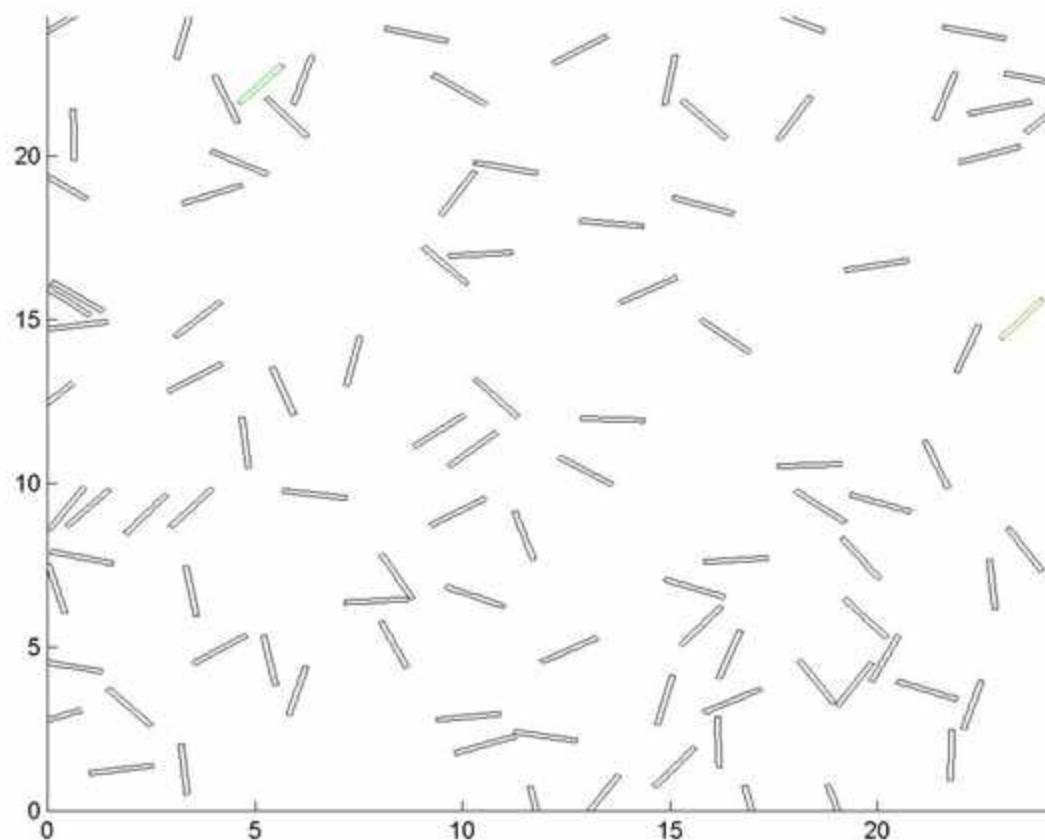


Interactions are due to overlapping of particles :

$$V(\vec{x}, \theta, \vec{x}', \theta') = \tilde{C} \left\{ \frac{1}{[\gamma - a(\vec{x}, \theta, \vec{x}', \theta')]^{\beta}} \frac{1}{\gamma^{\beta}} \right\}$$

# A specific example: collective motion in myxobacteria

- Putting the model in the computer- simulations w/ periodic boundary conditions!



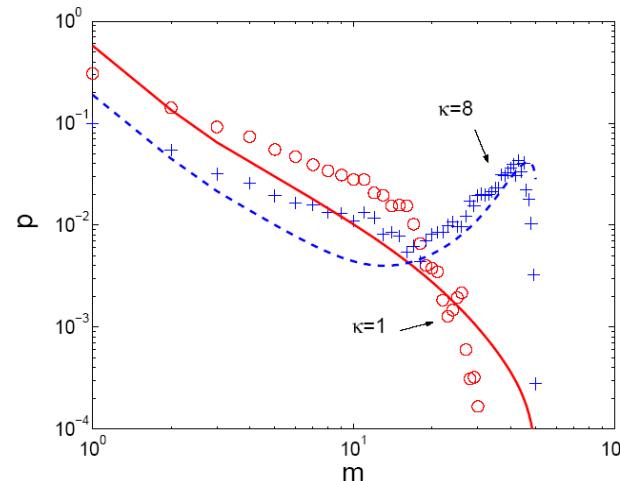
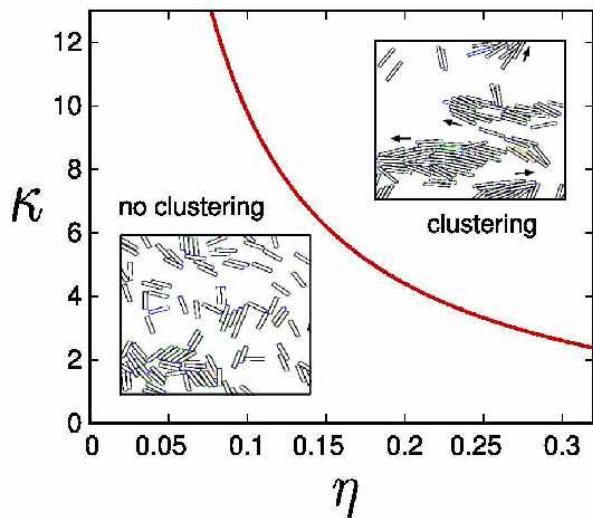
# A specific example: collective motion in myxobacteria

How can we characterize the macroscopic collective behavior of myxobacteria/self-propelled rods?



What can we measure here?

By looking at the clustering properties we can differentiate between individual and collective behavior



Peruani, Deutsch, and Bär, PRE (2006)

There is a dramatic change in the clustering properties of the system when either the density or the aspect ratio of the particles is changed!

# A specific example: collective motion in myxobacteria

The evolution equation for the cluster size distribution  $\rightarrow \{n_j(t)\}_{j=1}^{\infty}$

$$\dot{n}_1 = 2B_2 n_2 + \sum_{k=3}^N B_k n_k - \sum_{k=1}^{N-1} A_{k,1} n_k n_1,$$

$$\dot{n}_j = [B_{j+1} n_{j+1} - B_j n_j] - \sum_{k=1}^{N-j} A_{k,j} n_k n_j + \frac{1}{2} \sum_{k=1}^{j-1} A_{k,j-k} n_k n_{j-k} \quad \text{for } j = 2, \dots, N-1$$

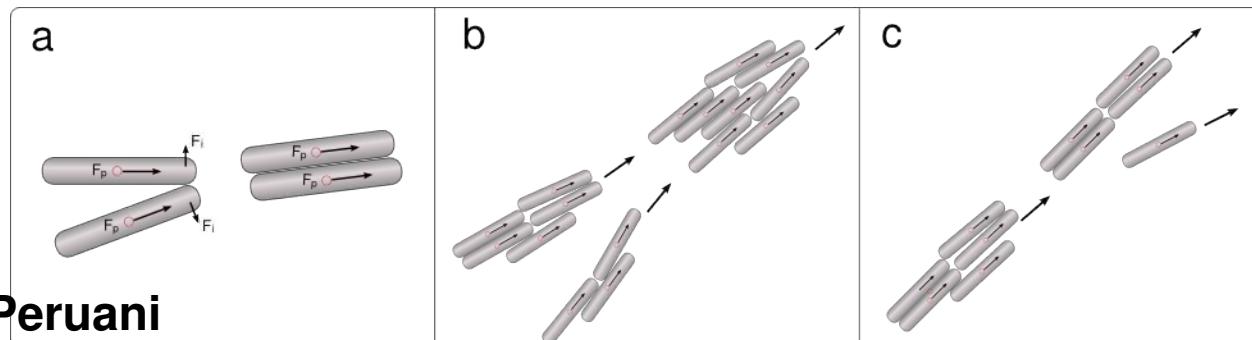
$$\dot{n}_N = -B_N n_N + \frac{1}{2} \sum_{k=1}^{N-1} A_{k,N-k} n_k n_{N-k},$$

collision kernel

$$A_{j,k} = (\tilde{v} \sigma_0 / A) (\sqrt{j} + \sqrt{k})$$

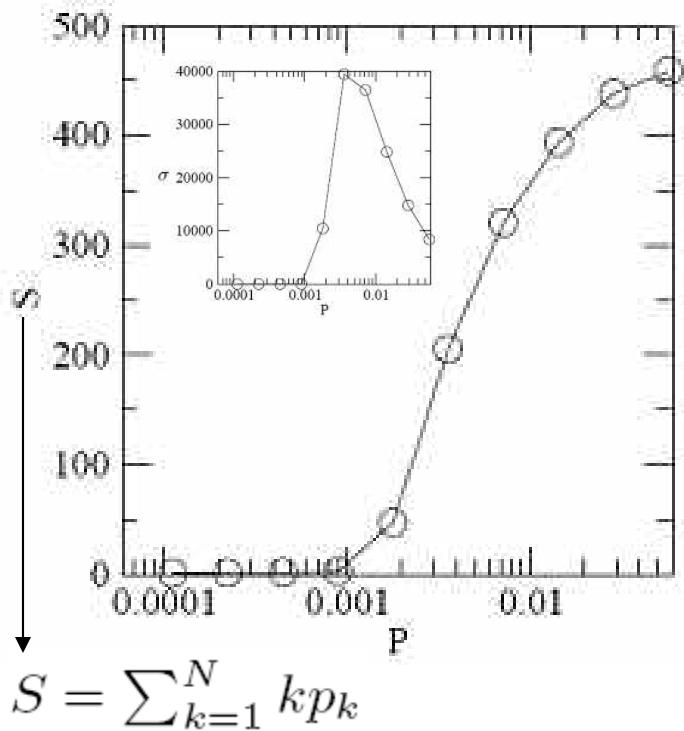
fragmentation kernel

$$B_j = (\tilde{v} / L) \sqrt{j}$$

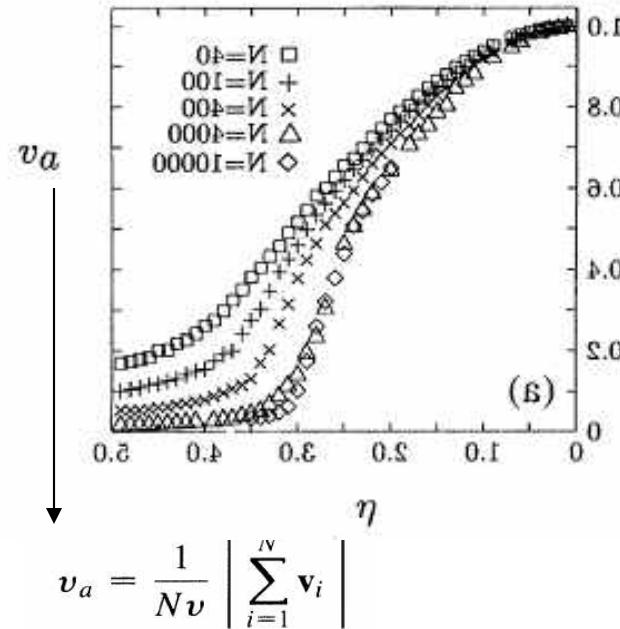


# A specific example: collective motion in myxobacteria

The system exhibits a phase transition to an aggregation phase



In Vicsek model...



Three types of distributions/phases:

- Exponential ( $P < P_c$ ) – [monodisperse phase w/ characteristic cluster size]
- Power-law ( $P = P_c$ ) – [at the transition point, scale-free distribution]
- Peak for large  $m$  ( $P > P_c$ ) – [aggregation phase!]

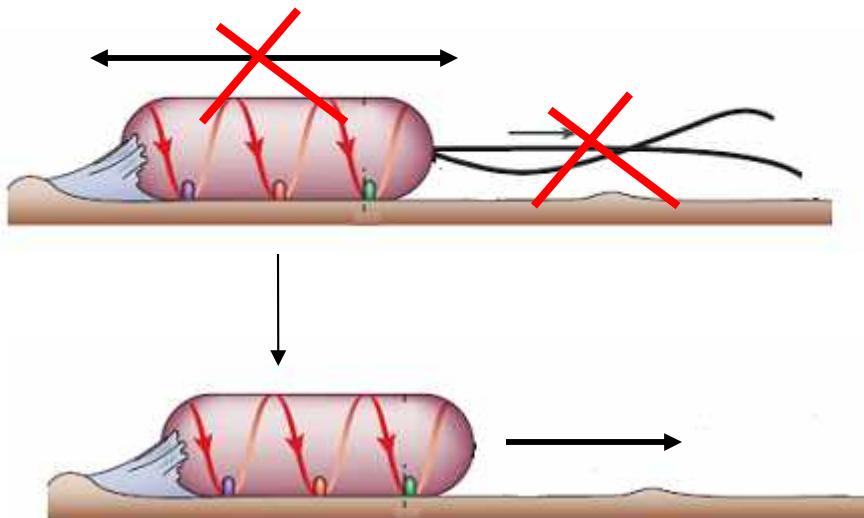
# A specific example: collective motion in myxobacteria

## What about collective motion and clustering in real myxobacteria ?

- Experiments with A+S-Frz- Myxococcus xanthus mutant

*In collaboration with:*

M. Bär, A. Deutsch, V. Jakovljevic, L. Søgaard-Andersen, and J. Starruß

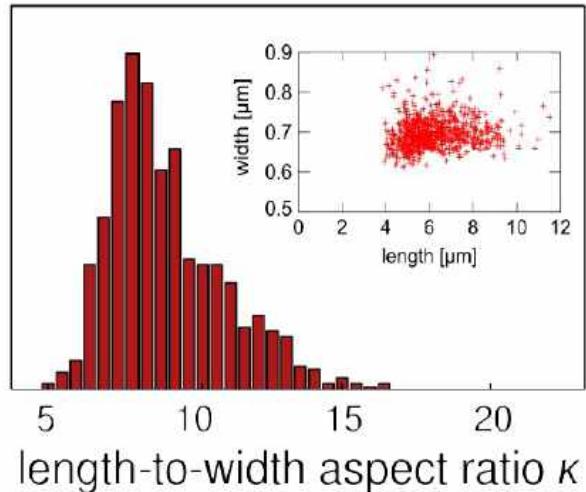


### Adventurous mutant:

- Cells do not reverse
- Social motility engine – off
- Advent. motility engine - on

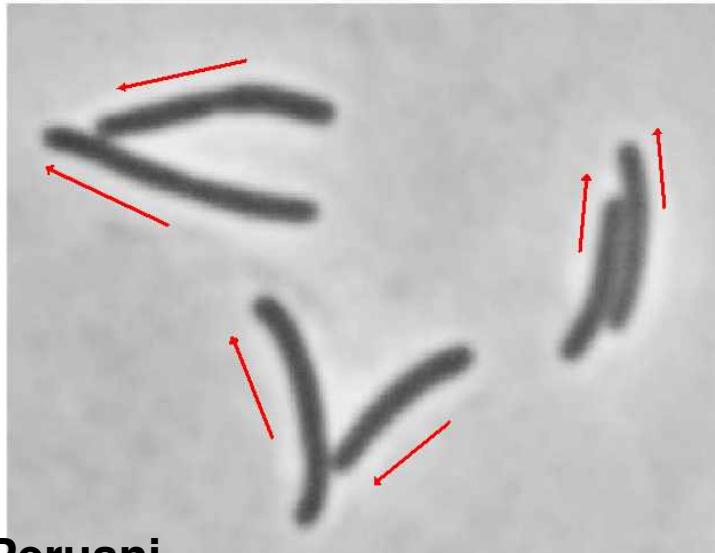
# A specific example: collective motion in myxobacteria

- Clustering in the A+S-Frz- mutant

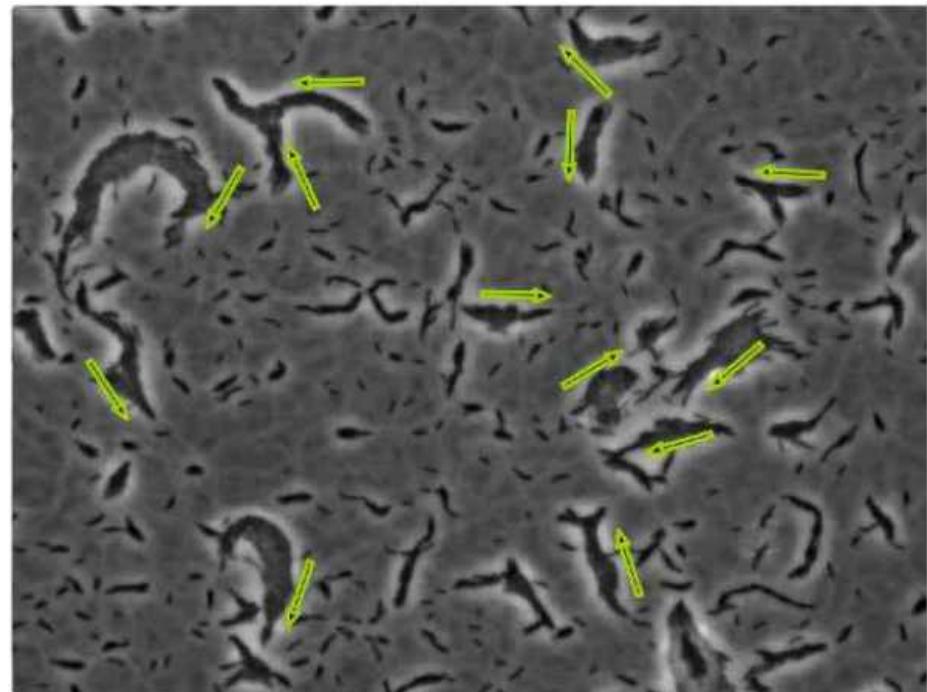


- Gliding speed =  $3.10 \pm 0.35 \mu\text{m}/\text{min}$
- $W=0.7 \mu\text{m}$ ,  $L=6.3 \mu\text{m}$ ,  $a=4.4 \mu\text{m}^2$
- $\kappa=8.9 \pm 1.95$

Cell collision leads to alignment:

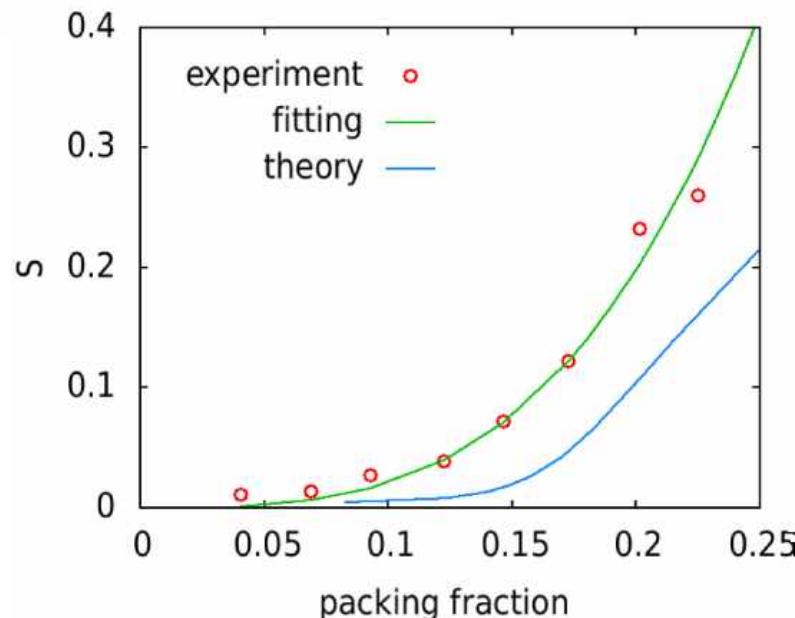
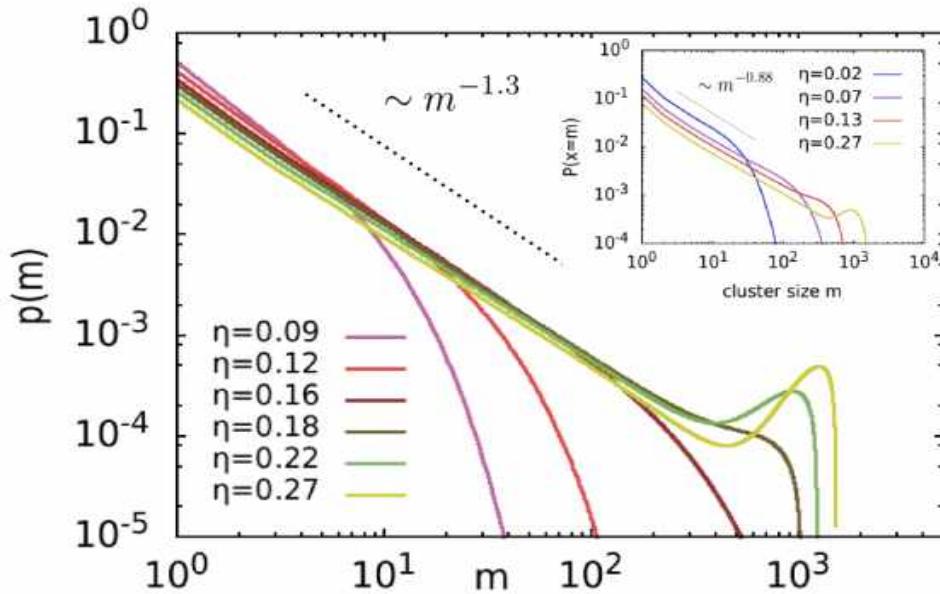
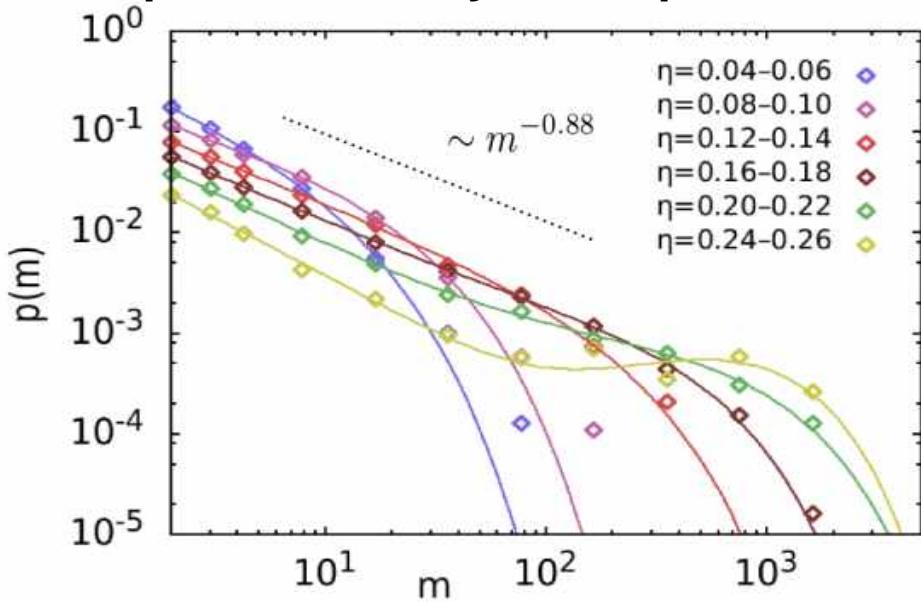


Moving clusters of bacteria are formed:

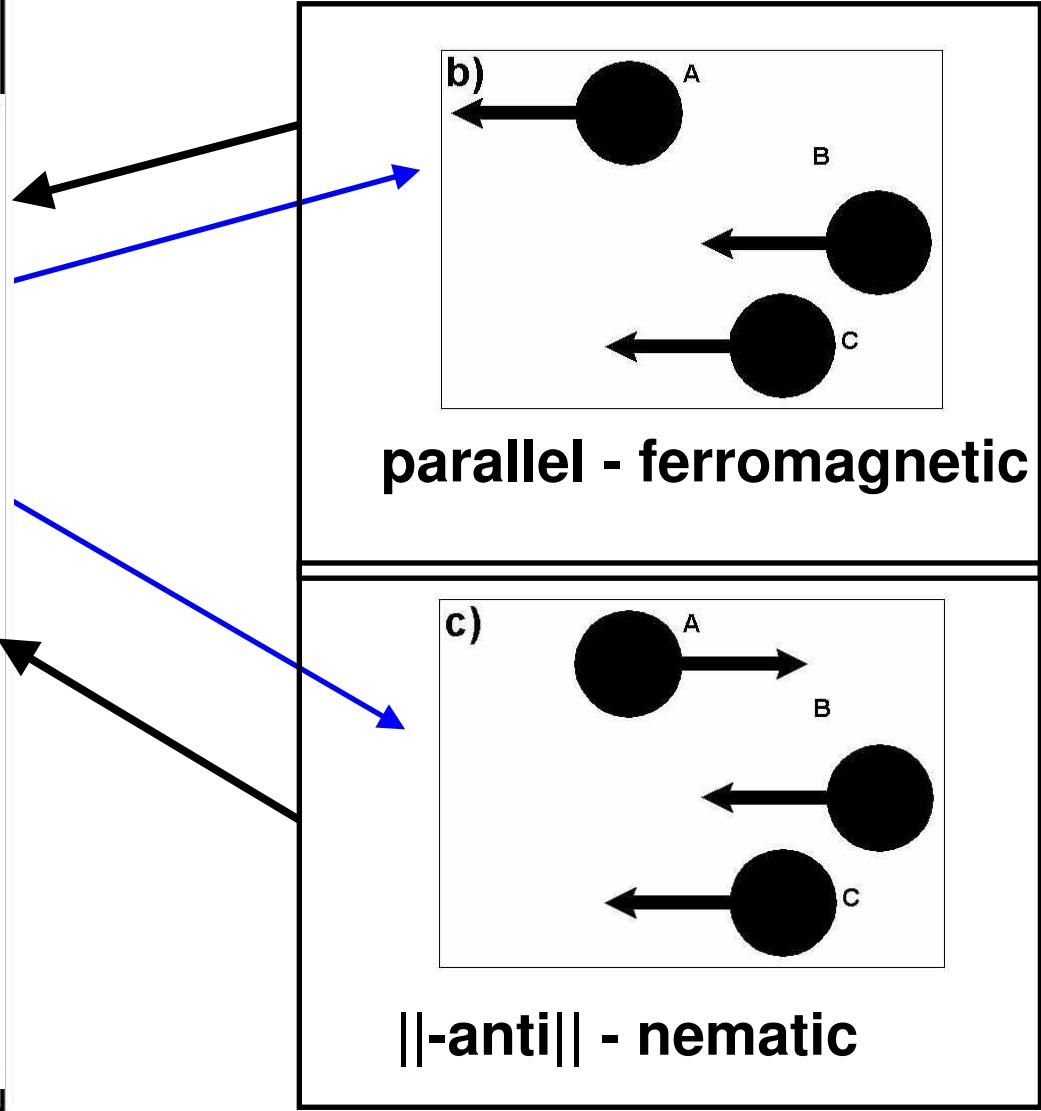
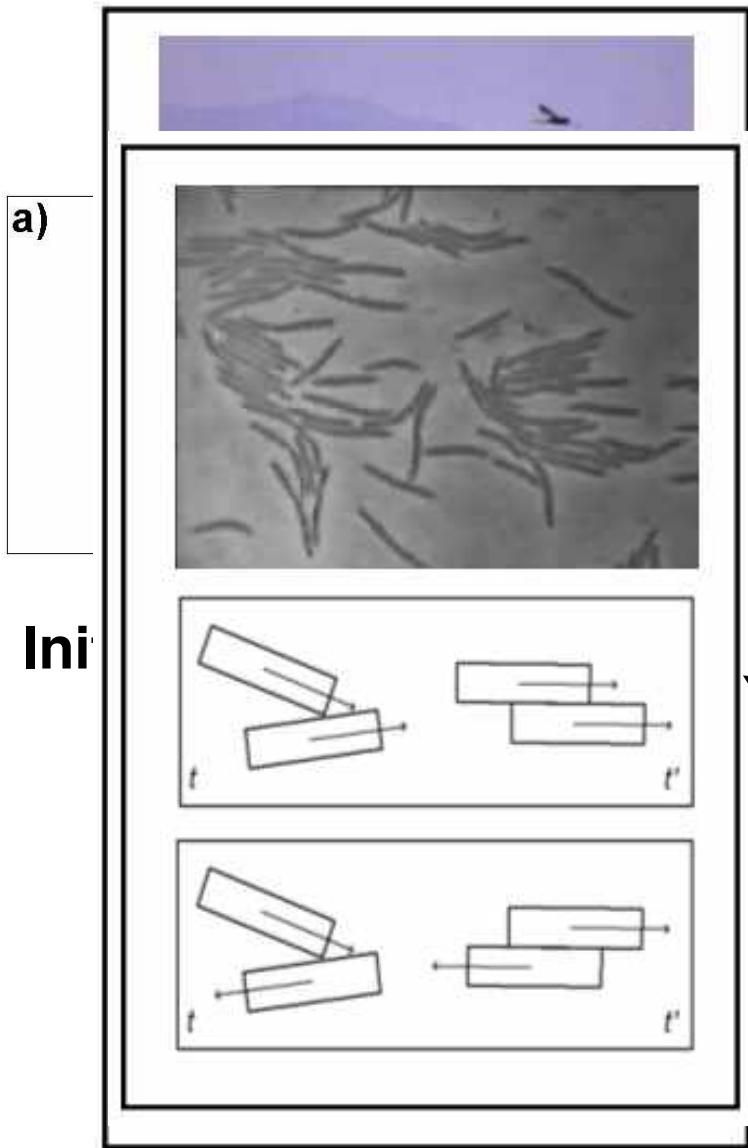


# A specific example: collective motion in myxobacteria

- Comparison: theory and experiments



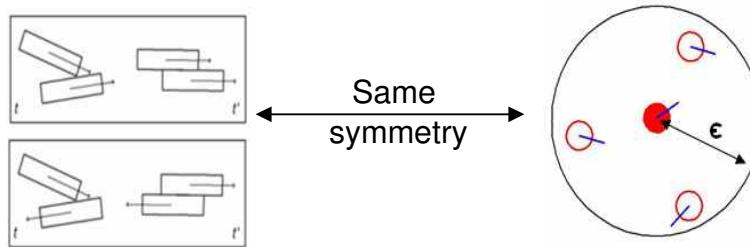
# Symmetries!



# Symmetries!

## A simple model for self-propelled rods (e.g., bacteria)

[F. Peruani, A. Deutsch, and M. Bär, Eur. Phys. J. Special Topics 157, 111 (2008)]



$$\mathbf{x}_i^{t+\Delta t} = \mathbf{x}_i^t + v_0 \mathbf{v}(\theta_i^t) \Delta t$$

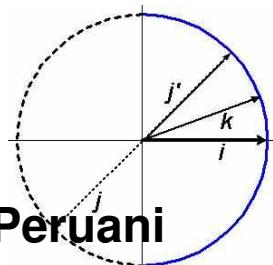
$$\theta_i^{t+\Delta t} = \arg \left( \sum_{|\mathbf{x}_i^t - \mathbf{x}_j^t| \leq \epsilon} \mathbf{f}(\mathbf{v}(\theta_j^t), \mathbf{v}(\theta_i^t)) \right) + \eta_i^t$$

Particles move in the direction given by:  
 $\mathbf{v}(\theta_i) = (\cos(\theta_i), \sin(\theta_i))$

Update of the moving direction

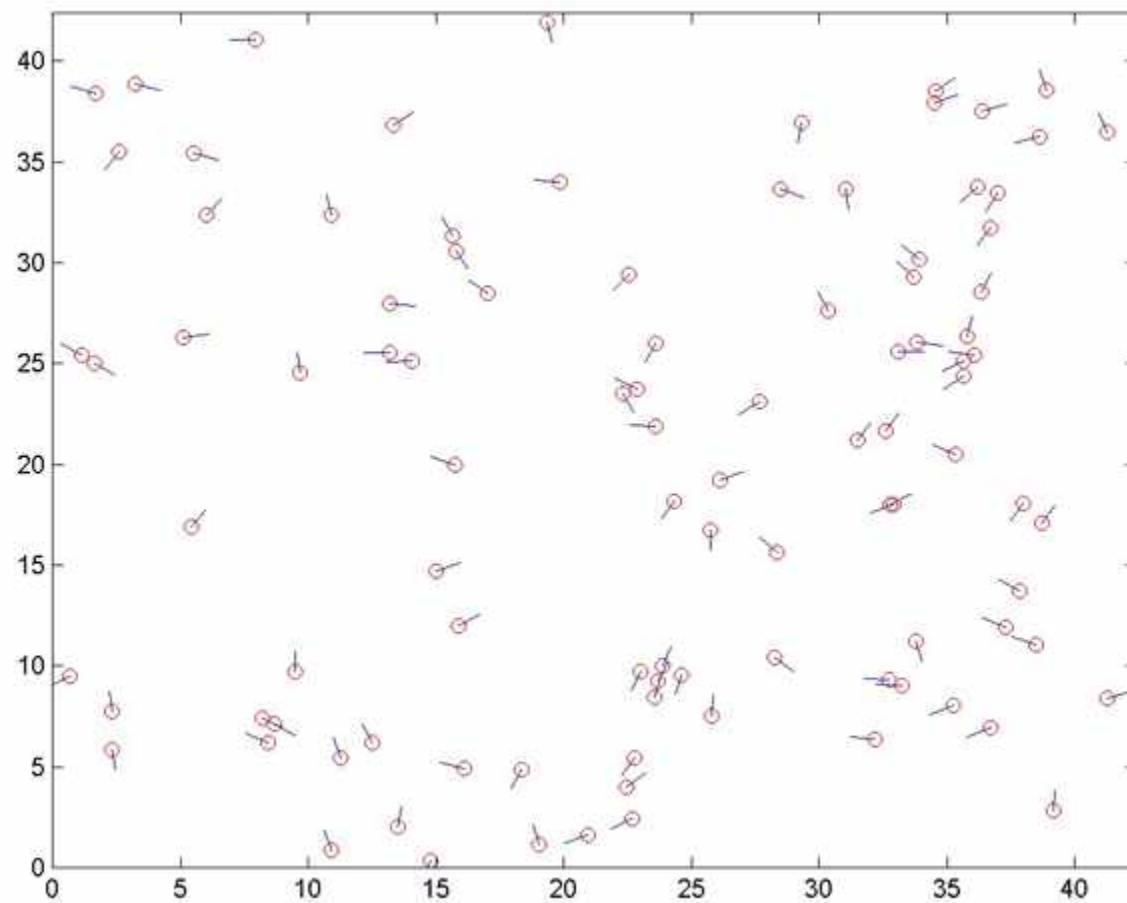
Alignment

Additive noise



$$f(\mathbf{v}_j, \mathbf{v}_i) = \begin{cases} \mathbf{v}_j & \text{if } \mathbf{v}_i \cdot \mathbf{v}_j \geq 0 \\ -\mathbf{v}_j & \text{if } \mathbf{v}_i \cdot \mathbf{v}_j < 0 \end{cases}$$

# Symmetries!



# Symmetries!

The symmetry of the alignment determines the type of macroscopic order

[F. Peruani, A. Deutsch, and M. Bär, Eur. Phys. J. Special Topics 157, 111 (2008)]

- A mean-field approach to understand collective motion

$$\partial_t C = D_\theta \partial_{\theta\theta} C + \gamma(\rho) \partial_\theta \left[ \underline{\partial_\theta \left( \int_0^{2\pi} d\theta' C(\theta', t) U(\theta, \theta') \right) C(\theta, t)} \right]$$

noise

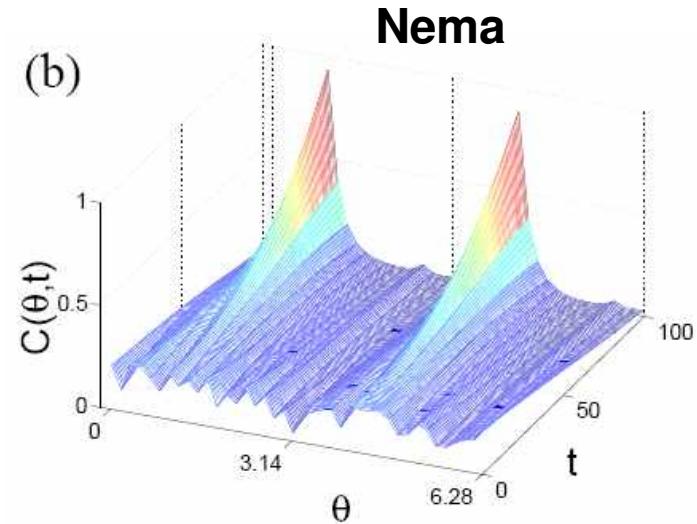
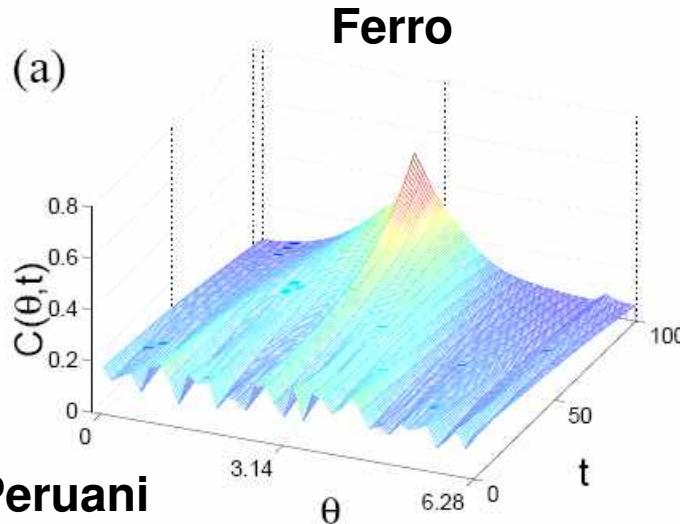
alignment

Ferromagnetic  
alignment

$$\rightarrow U(\theta, \theta') = -\cos(\theta - \theta')$$

Nematic alignment

$$\rightarrow U(\theta, \theta') = -\cos^2(\theta - \theta')$$



# Symmetries!

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[F. Peruani, A. Deutsch, and M. Bär, Eur. Phys. J. Special Topics 157, 111 (2008)]

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$$\partial_t C = D_\theta \partial_{\theta\theta} C + \gamma(\rho) \partial_\theta \left[ \underline{\partial_\theta \left( \int_0^{2\pi} d\theta' C(\theta', t) U(\theta, \theta') \right) C(\theta, t)} \right]$$

noise

alignment

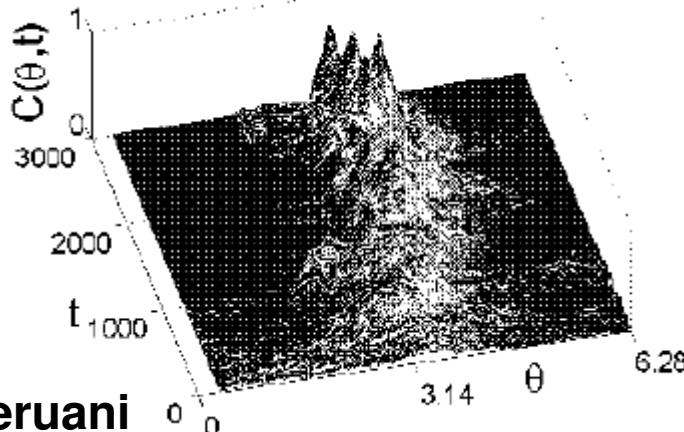
Ferromagnetic  
alignment

$$\rightarrow U(\theta, \theta') = -\cos(\theta - \theta')$$

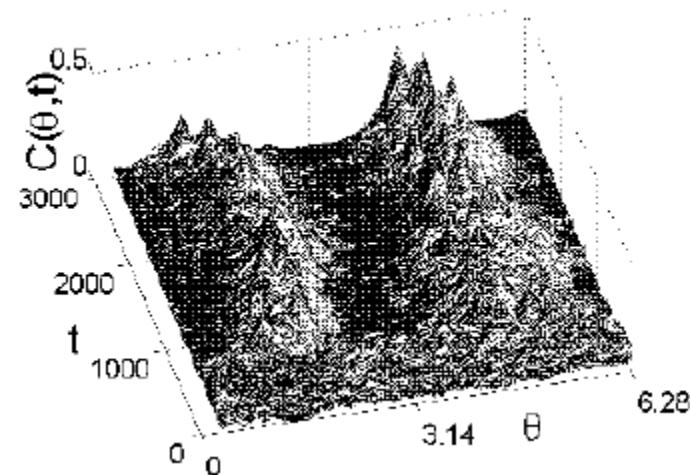
Nematic alignment  $\longrightarrow$

$$U(\theta, \theta') = -\cos^2(\theta - \theta')$$

Ferro



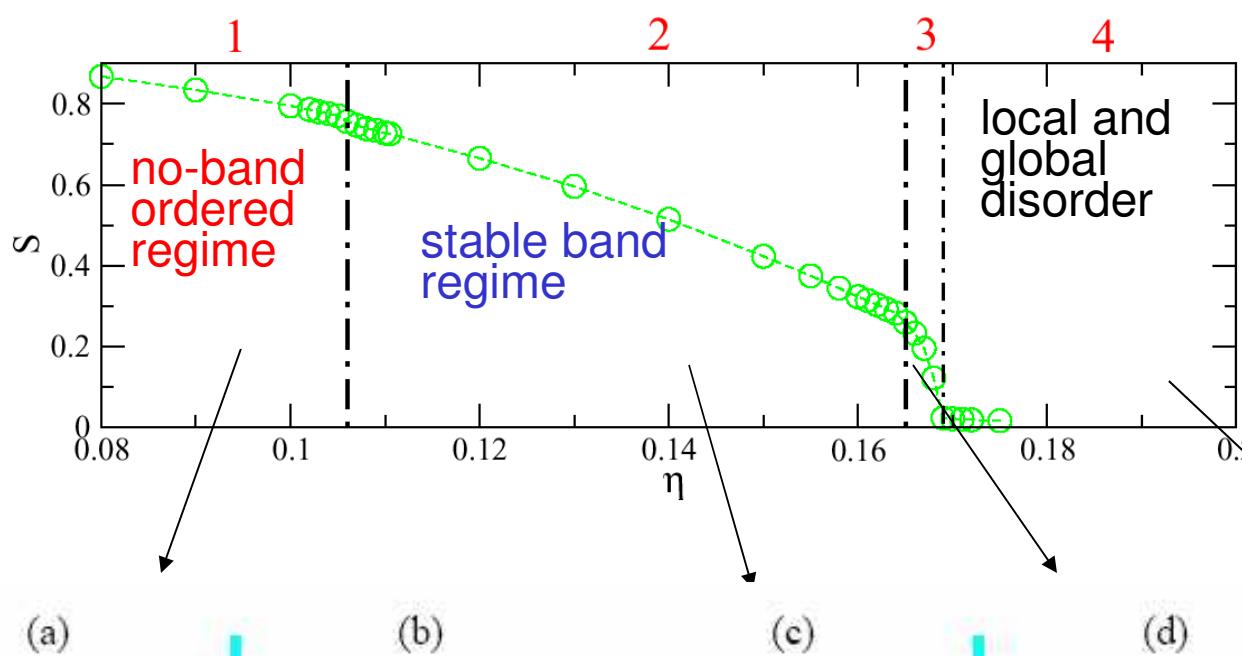
Nema



# Symmetries!

The symmetry of the alignment plays a crucial role in pattern formation

[F. Ginelli, F. Peruani, M. Bär, and H. Chaté, Phys. Rev. Lett. 104, 184502 (2010)]

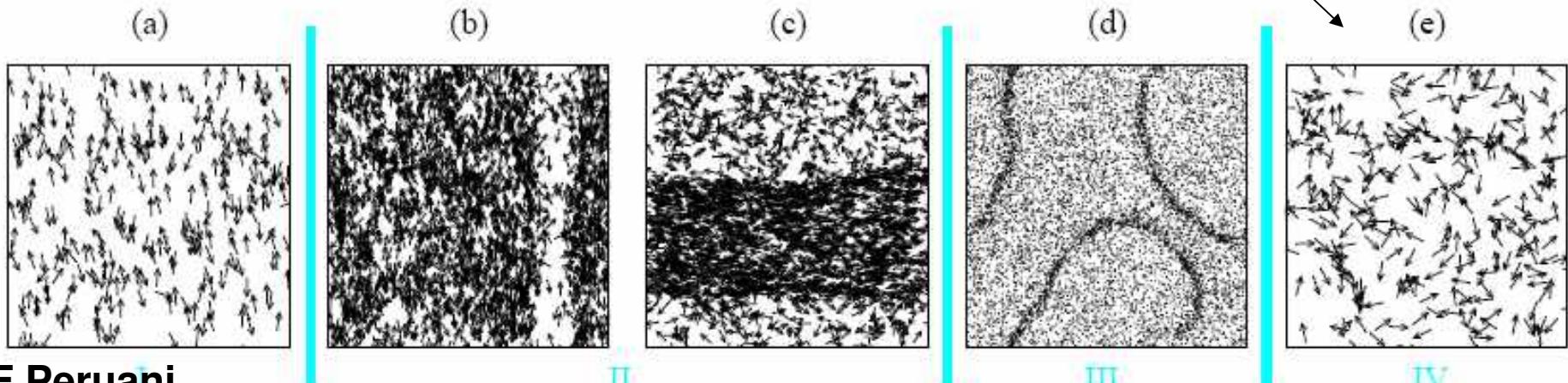


Nematic OP:

$$\langle S \rangle = \left\langle \left| \frac{1}{N} \sum_{k=1}^N e^{i 2 \theta_k} \right| \right\rangle_t$$

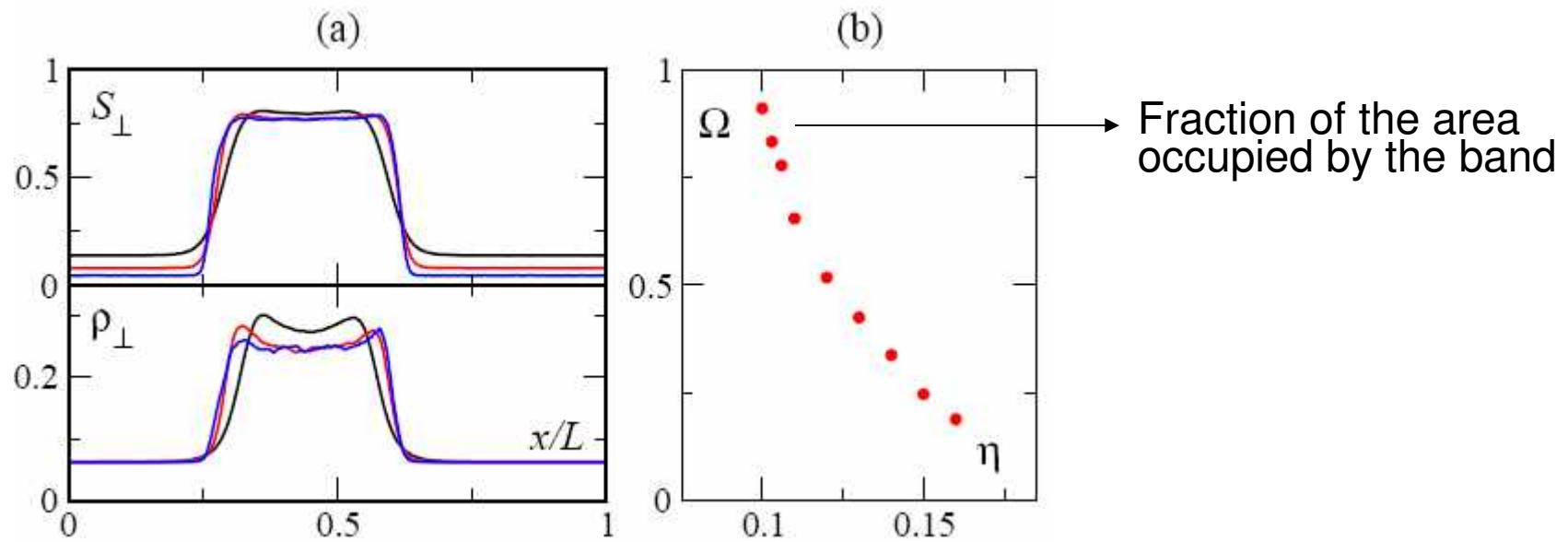
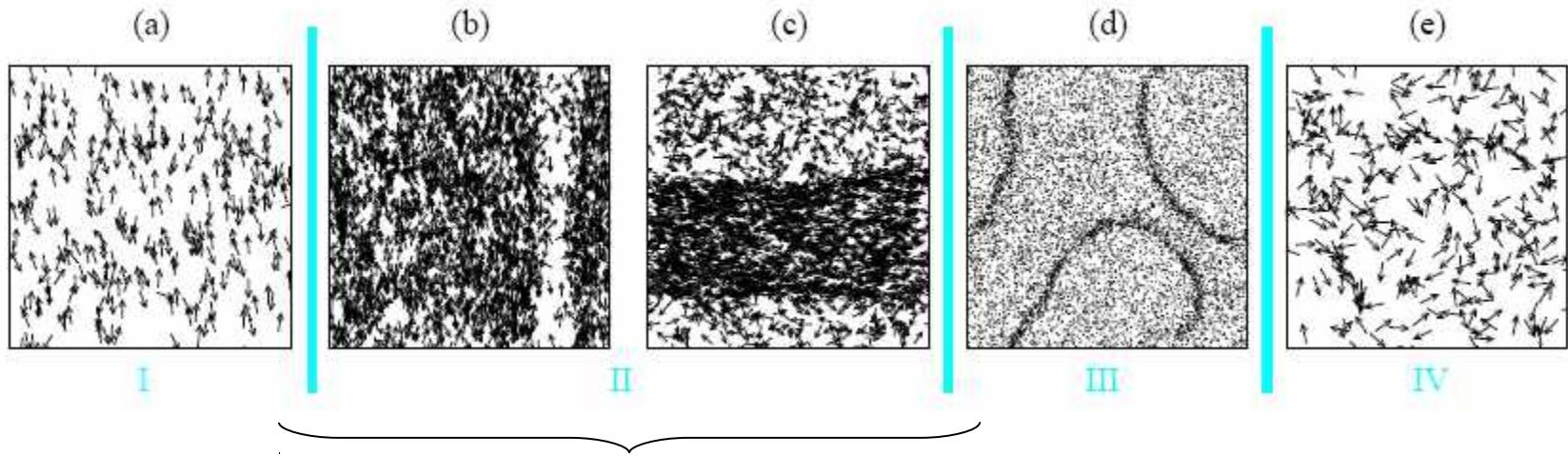
Ferromagnetic OP:

$$\langle \phi \rangle = \left\langle \left| \frac{1}{N} \sum_{k=1}^N e^{i \theta_k t} \right| \right\rangle_t$$

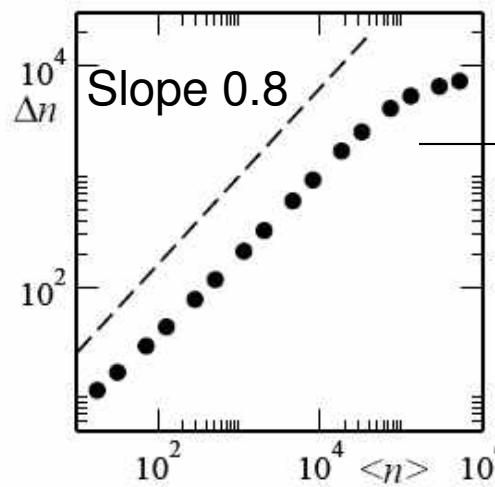
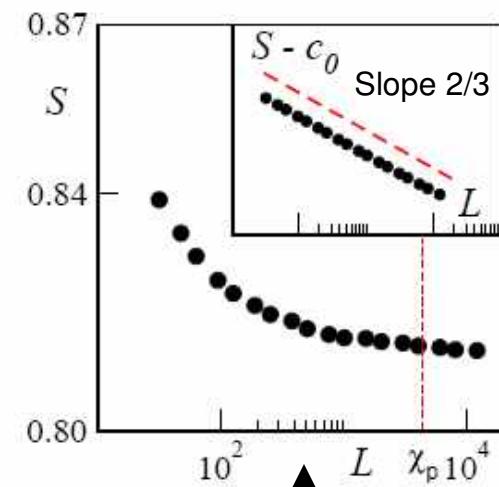


# Symmetries!

The symmetry of the alignment plays a crucial role in pattern formation



# Symmetries!

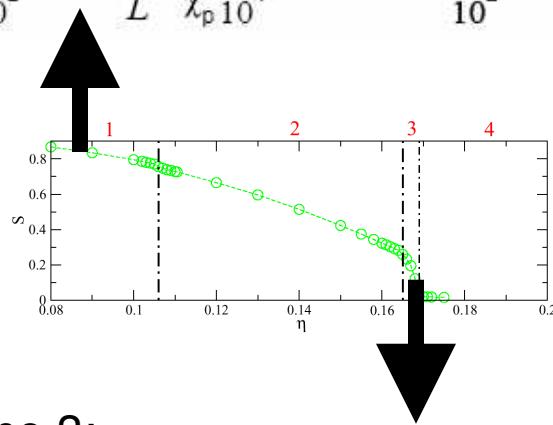
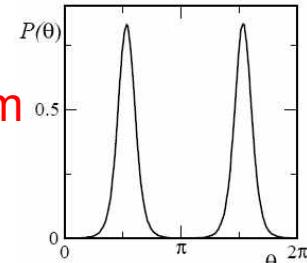


In regime 1:

- True Long-Range Order (LRO)

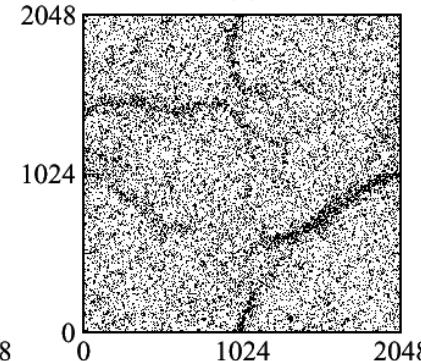
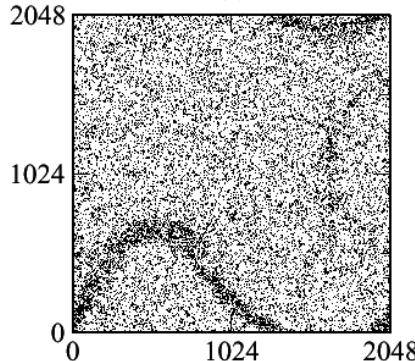
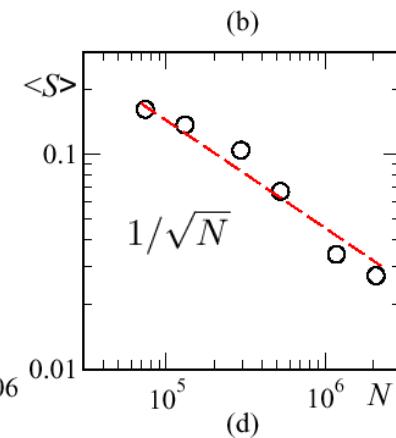
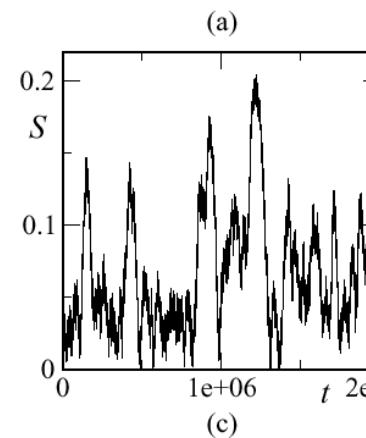
- Giant fluctuations

Mermin Wagner Theorem  
(for equilibrium syst.)  
does not allow for LRO!



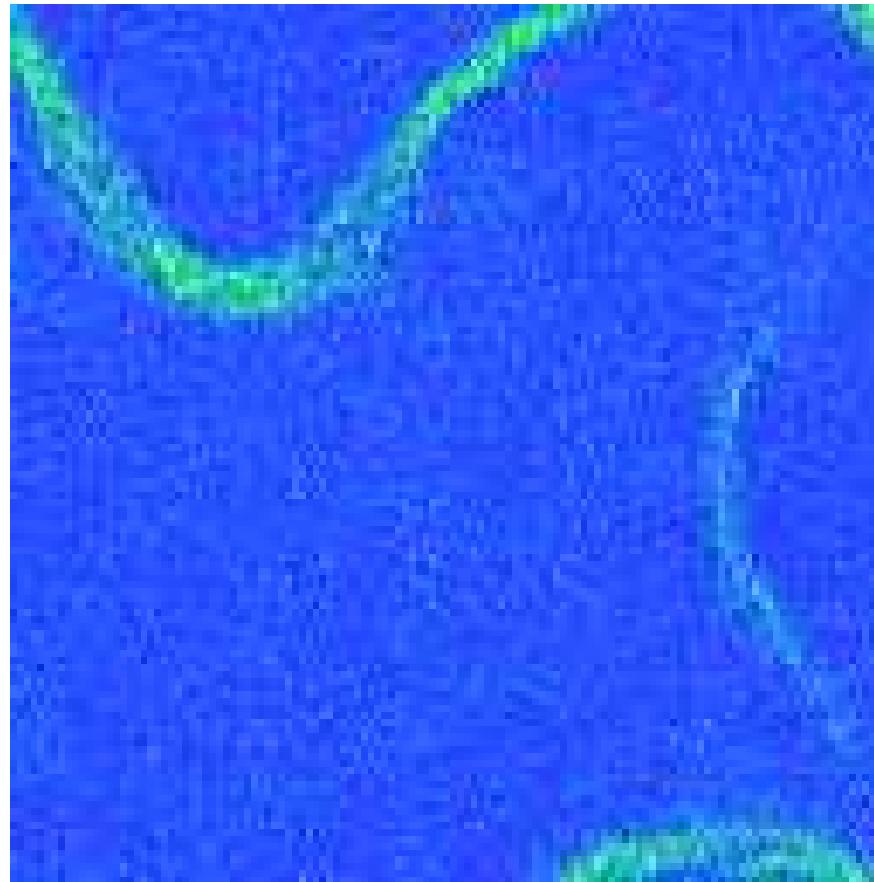
In regime 3:

- There is no LRO
- Unstable macroscopic structures (bands!)



# Symmetries!

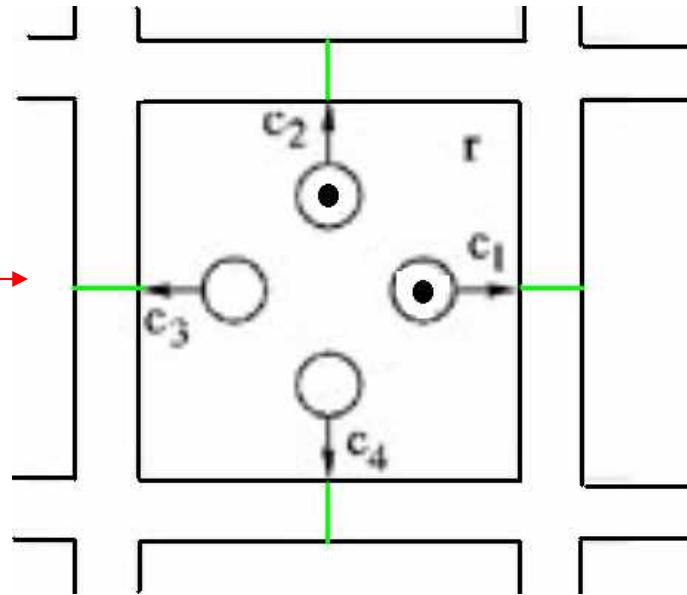
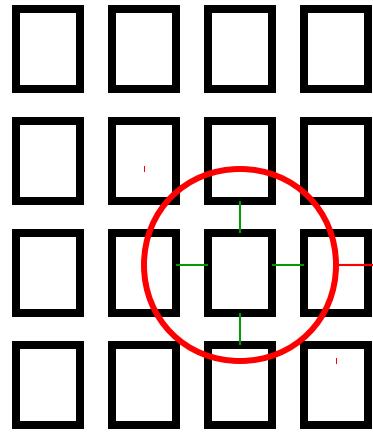
The symmetry of the alignment plays a crucial role in pattern formation



# Collective motion on the lattice

## Collective motion in a simple cellular automaton model

[H.J. Bussemaker, A. Deutsch, and E. Geigant, Phys. Rev. Lett. 78, 5018 (1997)]



### Rules:

- 1) Migration to next neighbor (according to channel direction)
- 2) Reorientation (i.e., velocity change) according to the following rule:

$$A[s \rightarrow \sigma \mid \mathbf{D}] = \frac{1}{Z} \delta[\rho(\sigma), \rho(s)] \exp[\beta \mathbf{D} \cdot \mathbf{J}(\sigma)]$$

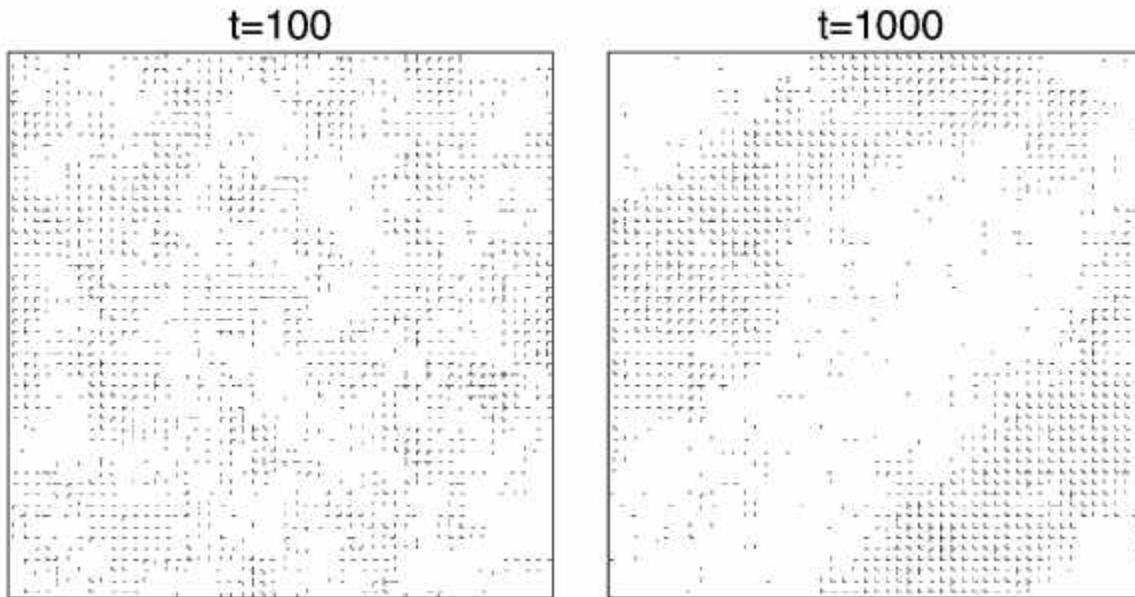
Mean local velocity:

Where  $Z$  defined such that  $\sum_{\sigma} A[s \rightarrow \sigma \mid \mathbf{D}] = 1$  and  $\mathbf{D}(\mathbf{r}, t) = \sum_{p=1}^4 \sum_{i=1}^4 \mathbf{c}_i s_i(\mathbf{r} + \mathbf{e}_p, t)$

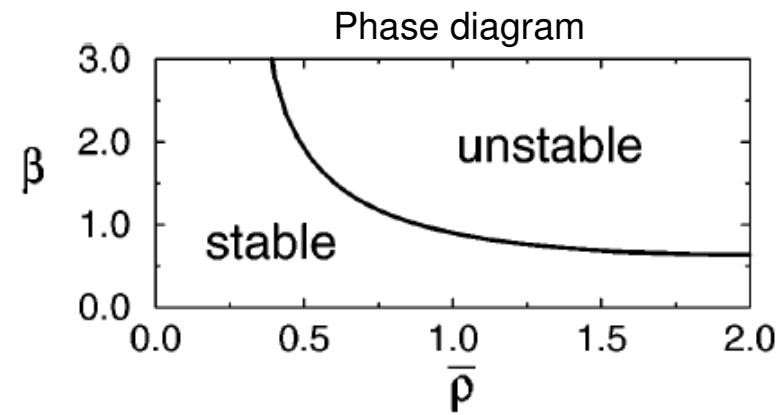
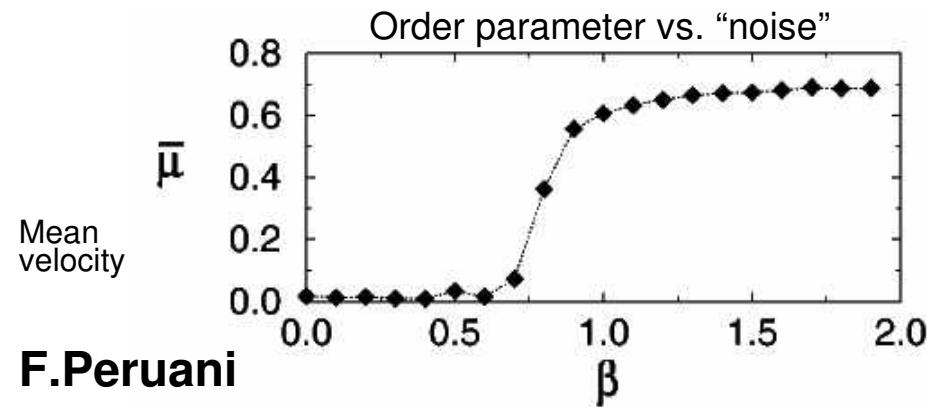
# Collective motion on the lattice

## Collective motion in a simple cellular automaton model

[H.J. Bussemaker, A. Deutsch, and E. Geigant, Phys. Rev. Lett. 78, 5018 (1997)]



Parameters:  $L=50$ ,  $\beta=1.5$ ,  $\rho=0.8$  ( $N \sim 2000$ )

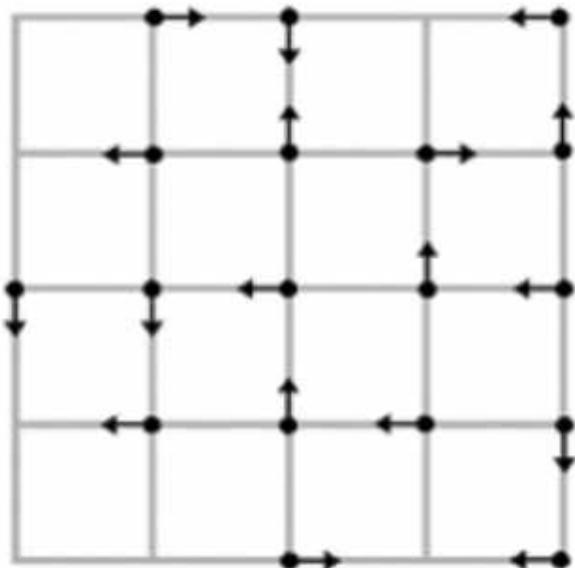


# Collective motion on the lattice

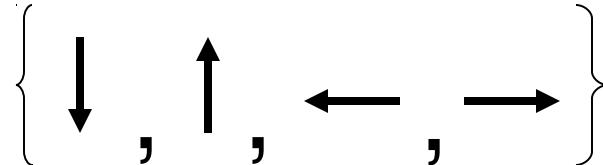
## First- and second order phase transition in a lattice model for swarming

In collaboration with:

T. Klauß, A. Deutsch, and A. Voß-Böhme



Possible velocities for a particle:

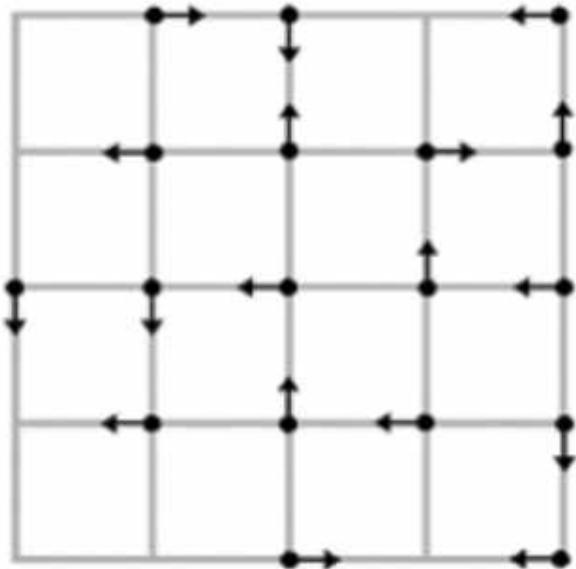


Possible states of a node:

- 1) Empty
- 2) Occupied

# Collective motion on the lattice

## First- and second order phase transition in a lattice model for swarming



The “stochastic rules”, defined in continuum time:

Each particle can perform two actions:

- 1) Migrate according to its velocity direction
- 2) Reorient its velocity direction

Associated to each action,  
there is a transition probability (per unit time):

$$T_M((x, v) \rightarrow (y = x + v, v)) = \begin{cases} m & \text{if node } y \text{ is empty} \\ 0 & \text{if node } y \text{ is occupied} \end{cases}$$

- 2) Reorient its velocity direction

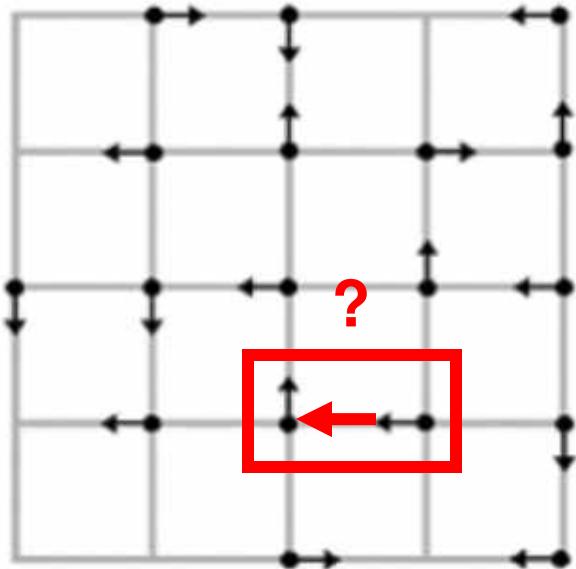
$$T_R((x, v) \rightarrow (x, w)) = \exp(g \sum_{y \in A(x)} \langle w | V(y) \rangle)$$

This defines a ferromagnetic alignment rule!

# Collective motion on the lattice

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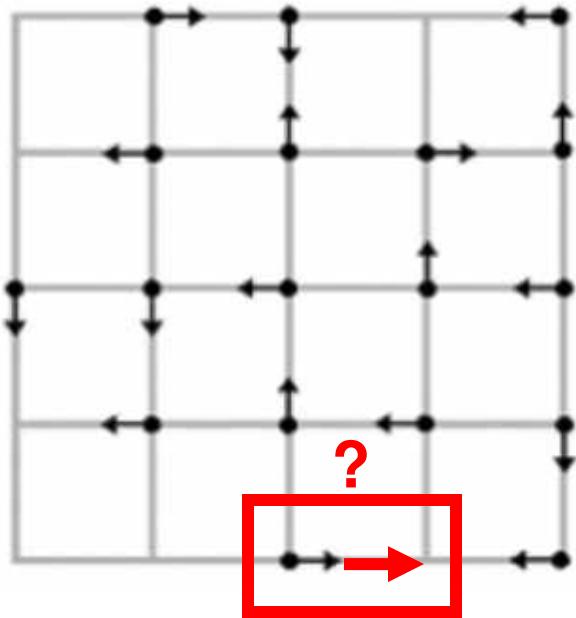
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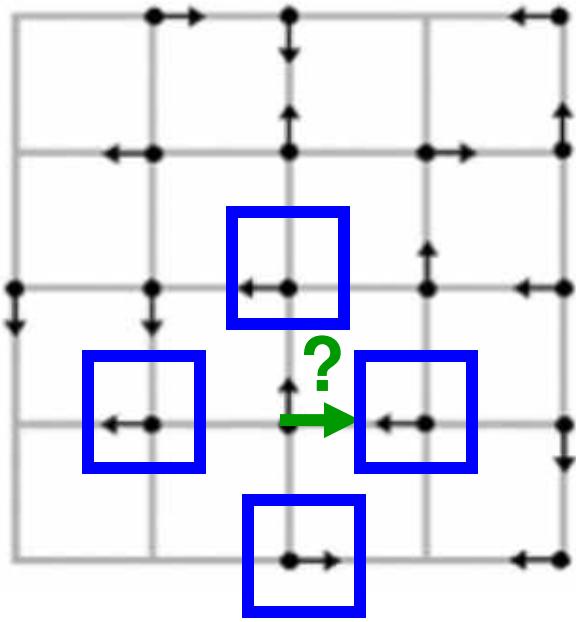
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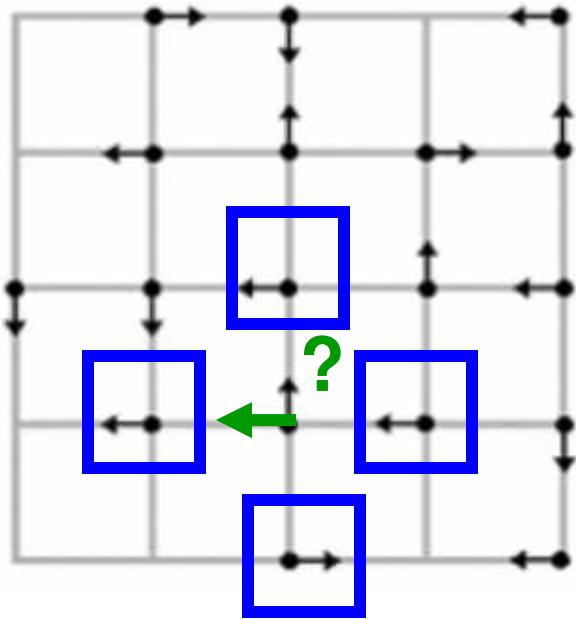
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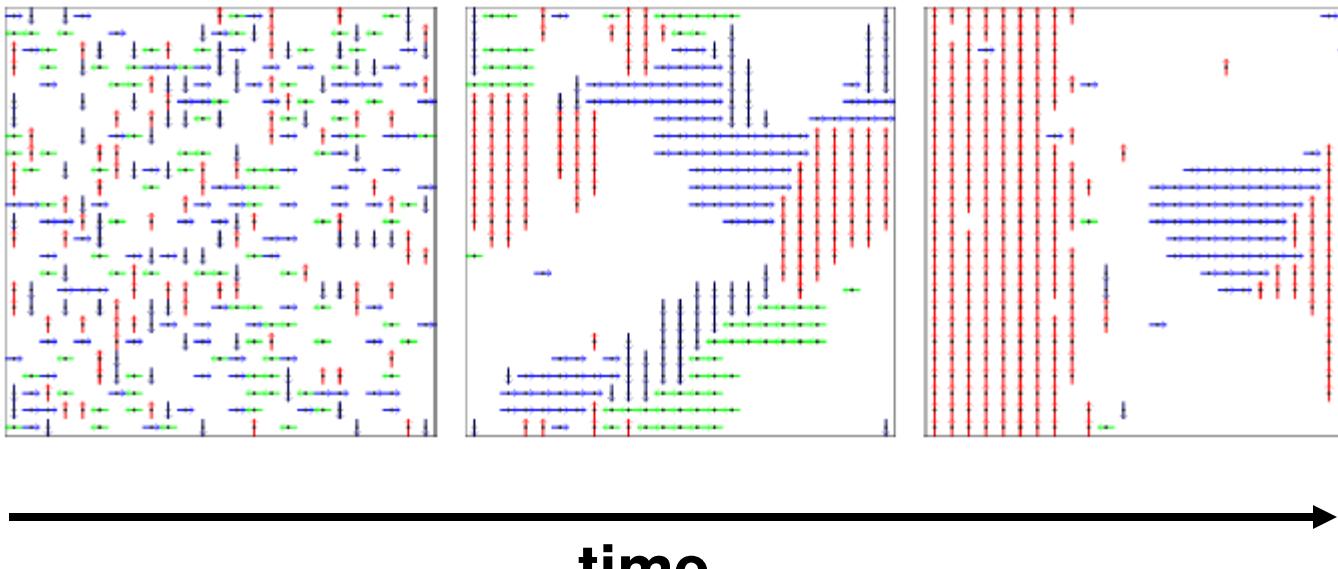
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# Collective motion on the lattice

## First- and second order phase transition in a lattice model for swarming

Results at “low” density

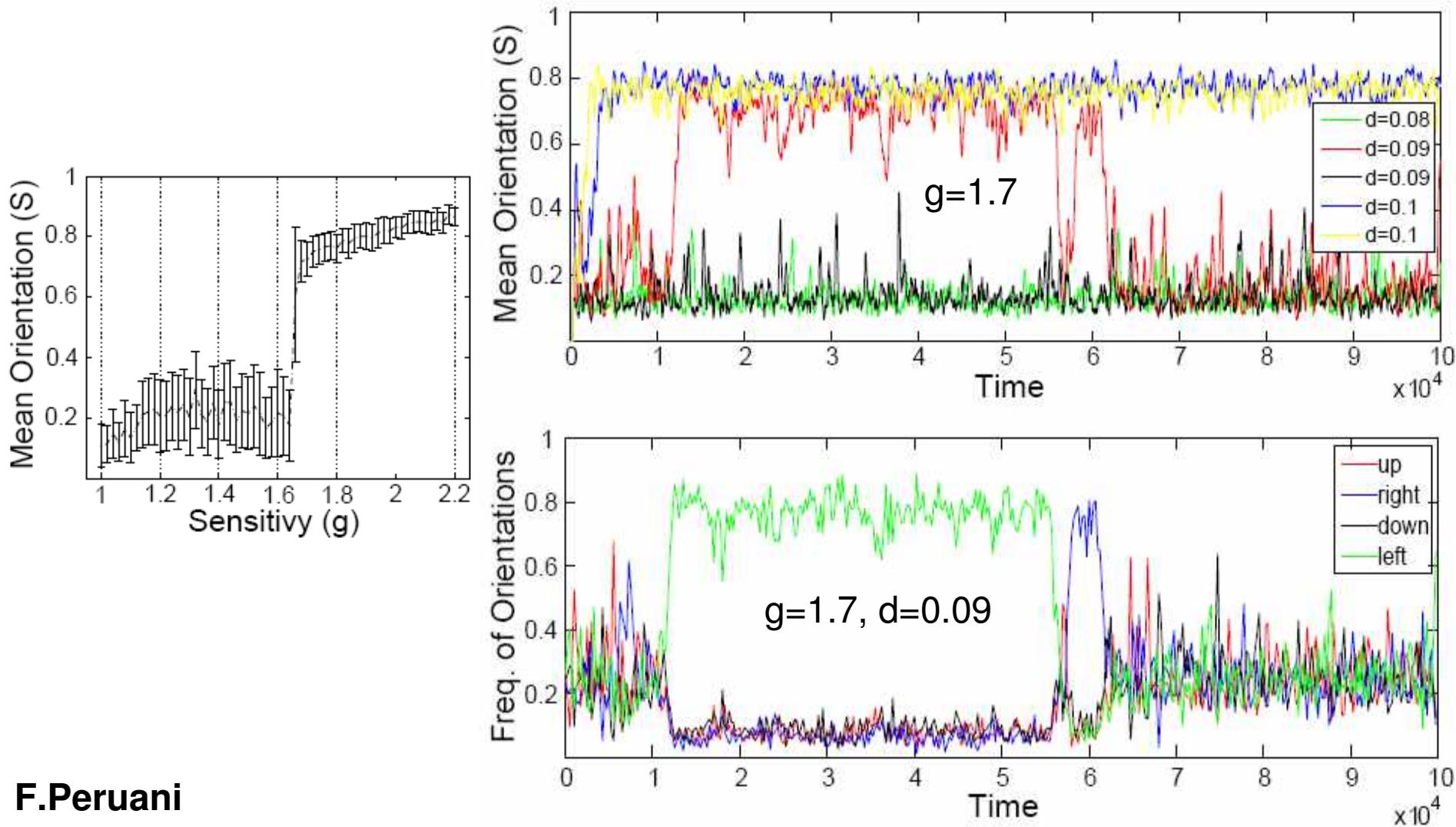


Parameters:  $L=25$ ,  $d=0.5$ ,  $g=1.7$ ,  $m=100$

# Collective motion on the lattice

## First- and second order phase transition in a lattice model for swarming

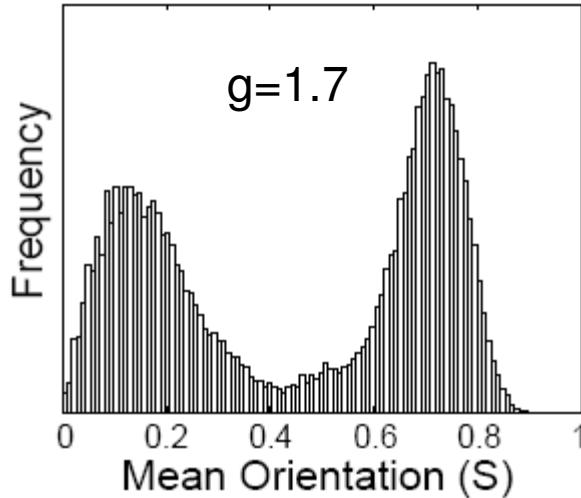
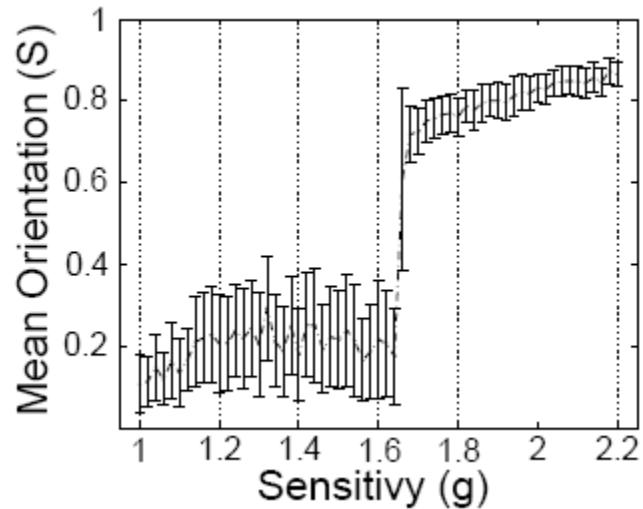
Results at “low” density ( $L=50, m=100$ )



# Collective motion on the lattice

## First- and second order phase transition in a lattice model for swarming

Results at “low” density



At “low” density the numerical evidence points towards a dynamical **first order** transition

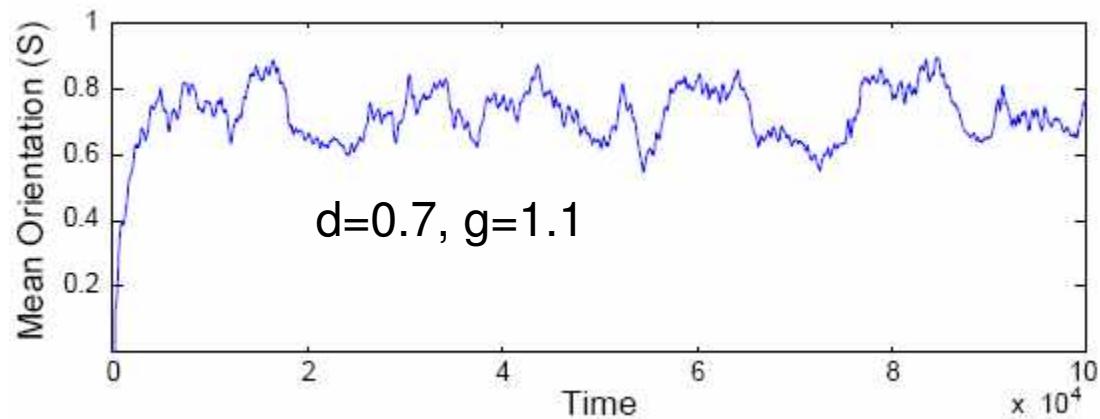
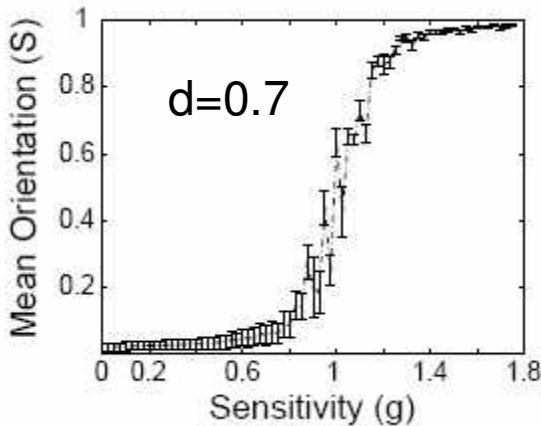
# Collective motion on the lattice

## First- and second order phase transition in a lattice model for swarming

Results at “high” (=full occupancy) density

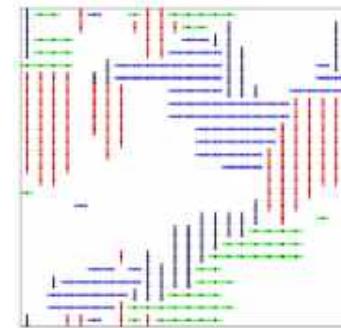
- The problem with full occupancy can be mapped to the 4-Potts model
- The 4-Potts model exhibits a second order phase transition
- Then, our model exhibits a **second order** transition at **d=1** !

Results at “intermediate” density

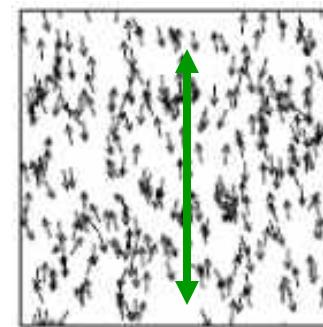
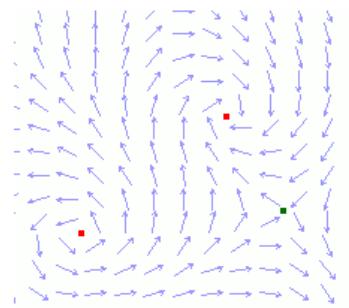


# Summary

- Models for collective motion can be either continuum (off-lattice) or discrete (on-lattice)

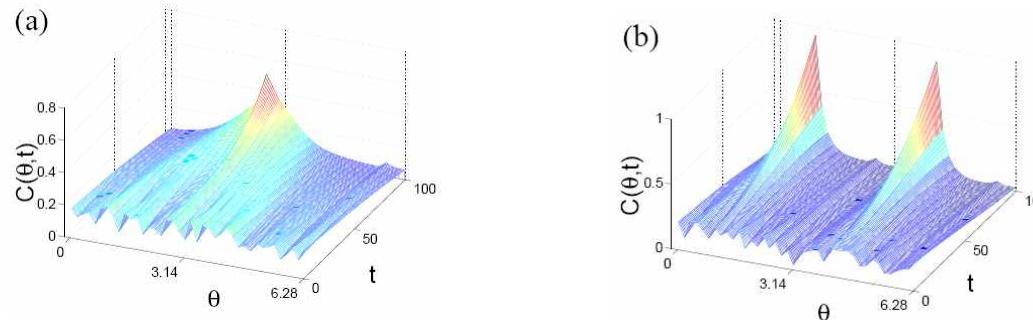


- Collective motion models are intrinsically (thermodynamically) non-equilibrium systems



# Summary

- The symmetry of the interaction determines the kind of order: either polar or nematic



- Collective motion can be characterized as a transition to aggregation (link to pattern formation)



- Collective motion modeling applies to **animal behavior** (fish, birds, sheeps, ants, etc) as well as **developmental biology** (aggregation patterns in bacteria, tissue formation, cancer, etc)



Thanks for you attention!

Some references:

FP, Deutsch, and Bär, PRE (2006)

FP, Morelli, PRL (2007)

FP, Deutsch, and Bär, EPJ-ST (2008)

Ginelli, FP, Bär, Chaté, PRL (2010)

FP, Nicola, Morelli, arXiv (2010)